

Enlightening the properties of the Higgs boson with the ATLAS Experiment using full Run-2 LHC data

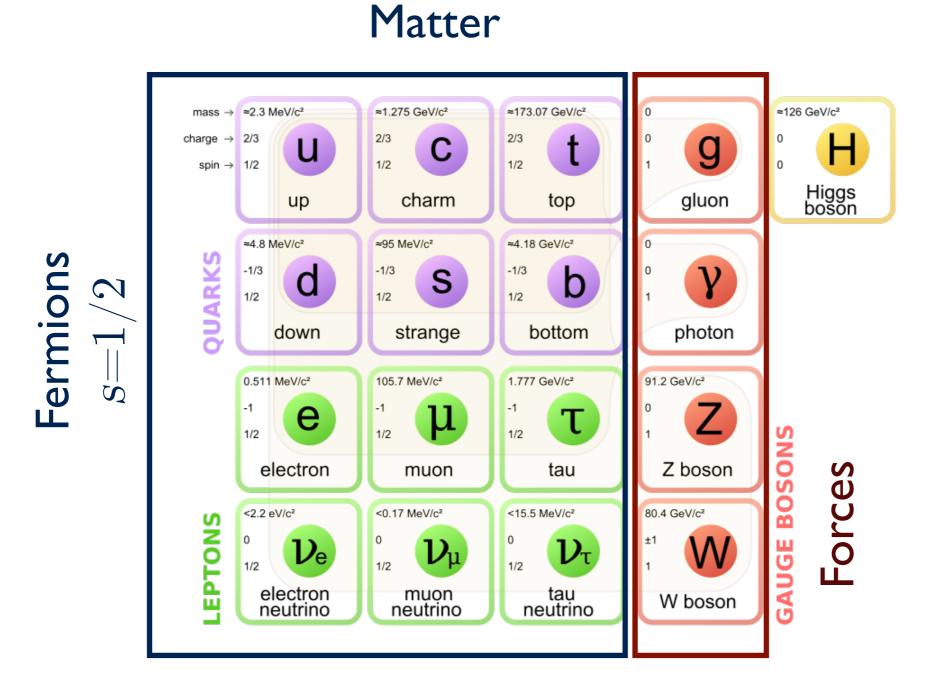
Gaetano Αθανάσιος Barone

Brookhaven National Laboratory



BNL, January 2021

Introduction



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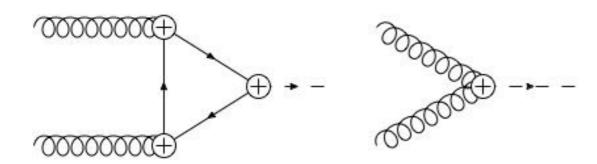
Fields



Higgs boson kinematics

Introduction

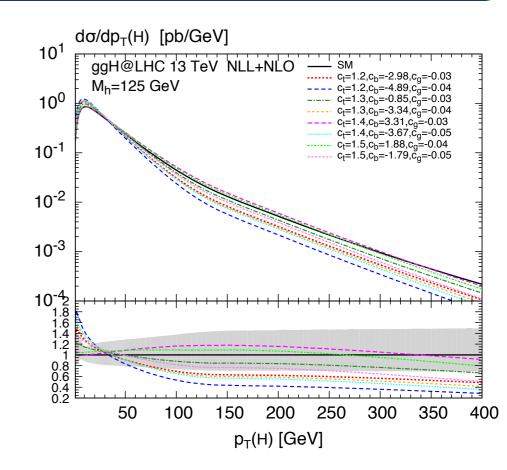
- Higgs boson kinematics:
 - ▶ $p_{T,4\ell}$: Lagrangian structure of H interactions.

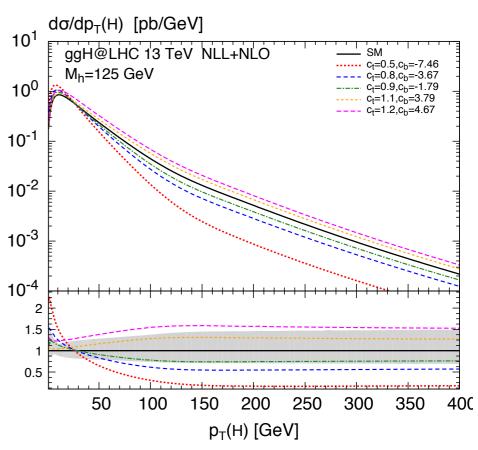


- Sensitivity to phenomena resonant at higher energies
 - → changes in observables at lower energies.

$$\begin{array}{l} \frac{c_1}{\Lambda^2}\,\mathcal{O}_1 \to \frac{\alpha_{\rm S}}{\pi v}c_g h G^a_{\mu\nu}G^{a,\mu\nu}\,, \, \Big\} \quad \text{cg: } ggH \text{ contact interaction} \\ \frac{c_2}{\Lambda^2}\,\mathcal{O}_2 \to \frac{m_t}{v}c_t h \bar{t}t\,, \\ \frac{c_3}{\Lambda^2}\,\mathcal{O}_3 \to \frac{m_b}{v}c_b h \bar{b}b\,, \, \Big\} \quad \text{ct: } t \text{ and } b \text{ Yukawa couplings} \\ \frac{c_4}{\Lambda^2}\,\mathcal{O}_4 \to c_{tg}\frac{g_S m_t}{2v^3}(v+h)G^a_{\mu\nu}(\bar{t}_L\sigma^{\mu\nu}T^at_R+h.c) \end{array}$$

 c_{tg} : dipole-moment, g-t interaction

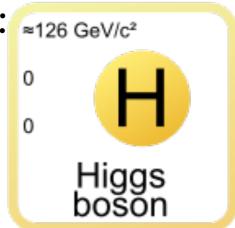


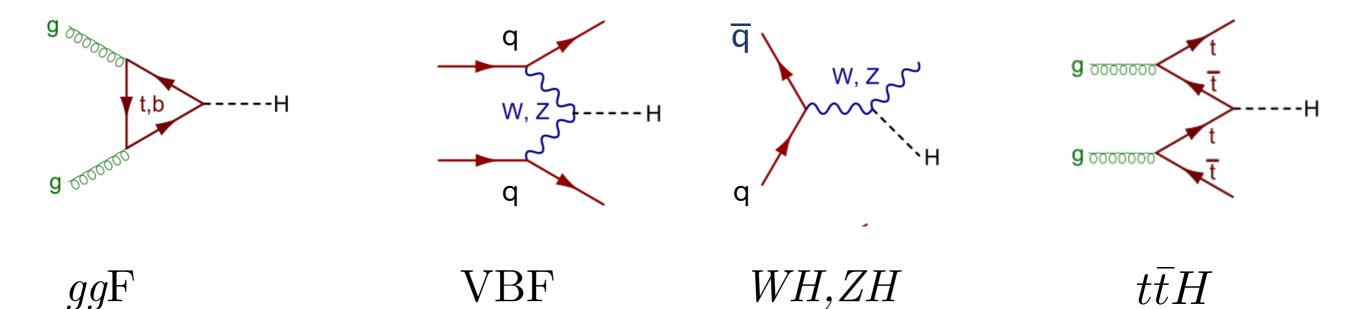




Introduction

- Higgs boson (H) main production in proton-proton collisions: $_{\sim 126 \text{ GeV/c}^2}$
 - ▶ Predominant production gluon-gluon fusion (87%) and VBF (6.8)
 - W,Z associated production (4%) and $t\overline{t}H$ (<1%)

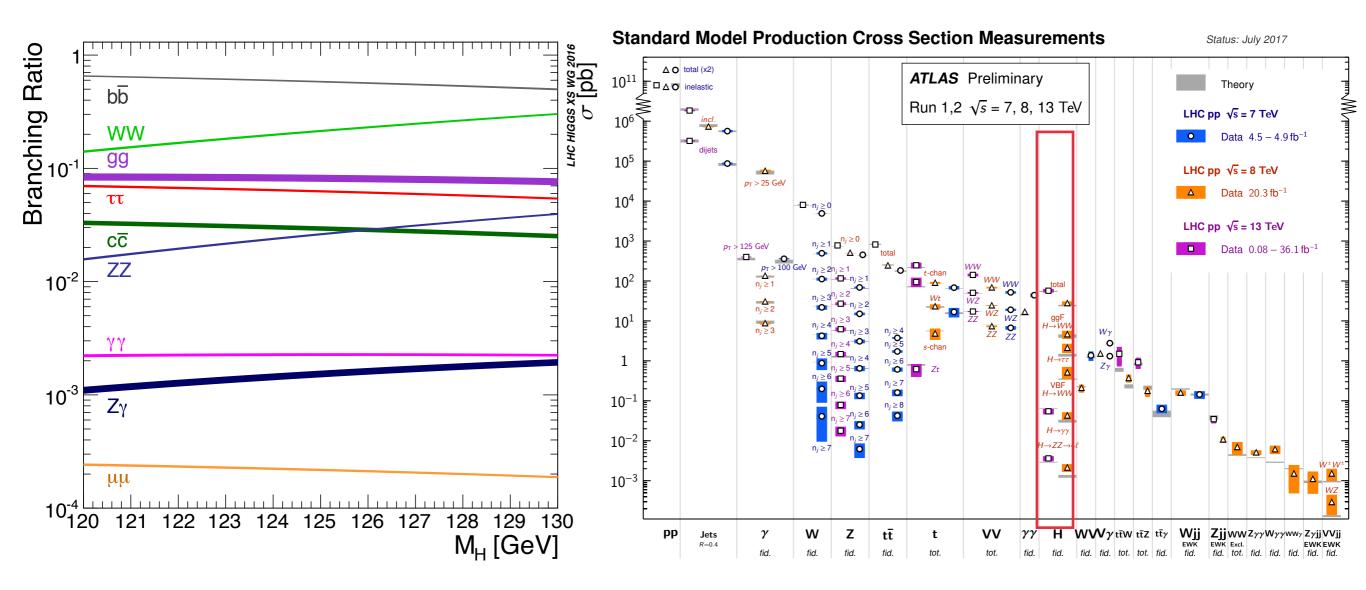




- Experimentally challenging final states
 - In association with additional jets, same final state of many other processes
 - forward jets with large rapidity gap, small rates



Introduction



- Channels considered in this talk :
 - (i) $H \rightarrow Dibosons (ZZ^* \rightarrow 4\boldsymbol{\ell}, WW^* \rightarrow \boldsymbol{\ell} \overline{\nu} \overline{\boldsymbol{\ell}} \nu)$, and $\gamma \gamma$
 - (ii) $H \rightarrow \text{light leptons } (ee, \mu \overline{\mu})$

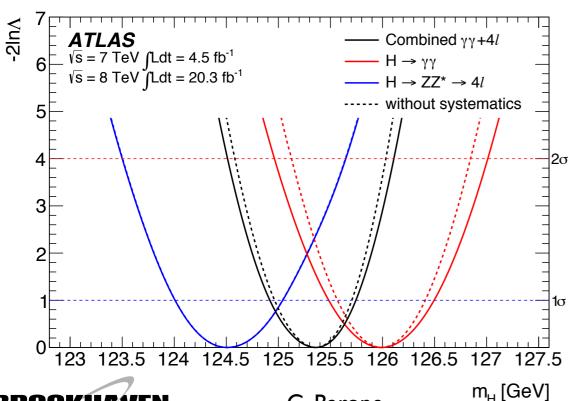


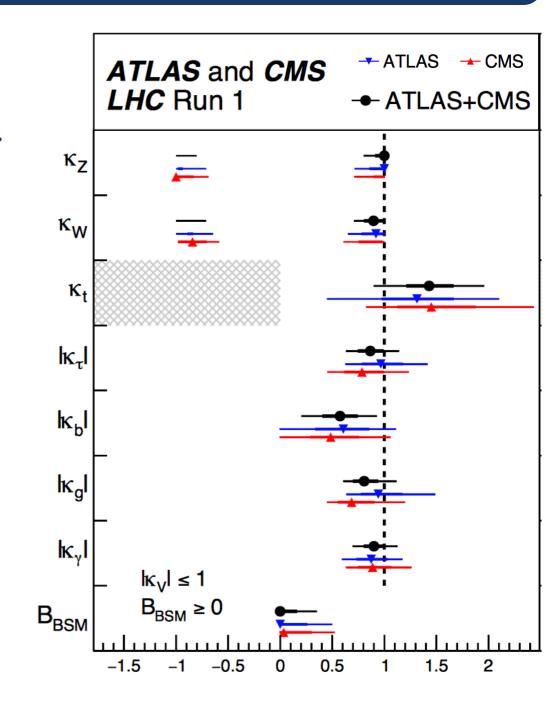
Run I Legacy

- Run-I featured in primis the discovery
 - First properties measurements
 - Programme largely limited by statistical accuracy.

Properties:

- ▶ ATLAS precision in m_H of 0.33%:
- ▶ Couplings measured to 10% to 25% precision
- ▶ $H \rightarrow \text{inv.}$ constrained to < 30%
- First studies of $J^{\rm PC}=0^{++}, {
 m (indirect)}$ width $\Gamma_{H}\!\!<14.4~{
 m MeV}~(15.2~{
 m MeV})$



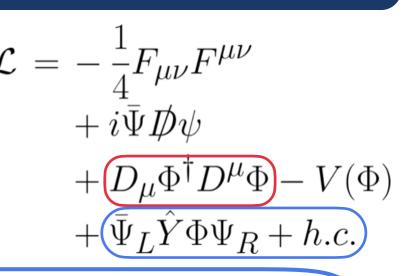


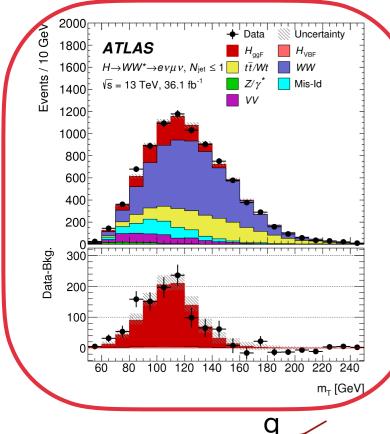
Channel	Mass measurement [GeV]	
$H \rightarrow \gamma \gamma$	$125.98 \pm 0.42 \text{ (stat)} \pm 0.28 \text{ (syst)} = 125.98 \pm 0.50$	
$H \rightarrow ZZ^* \rightarrow 4\ell$	$124.51 \pm 0.52 \text{ (stat)} \pm 0.06 \text{ (syst)} = 124.51 \pm 0.52$	
Combined	$125.36 \pm 0.37 \text{ (stat)} \pm 0.18 \text{ (syst)} = 125.36 \pm 0.41$	

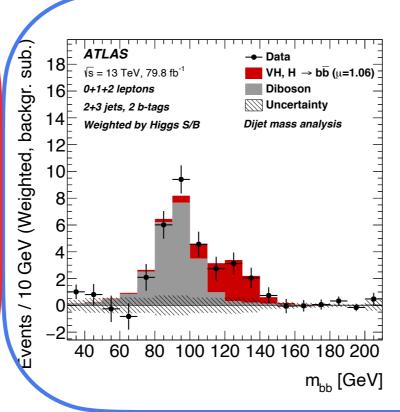
Overview

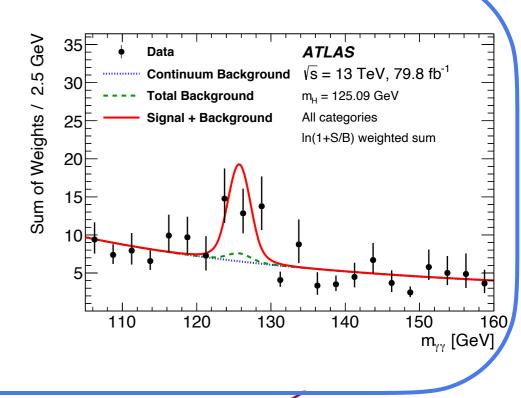
ATLAS collected 139 fb⁻¹ in Run 2

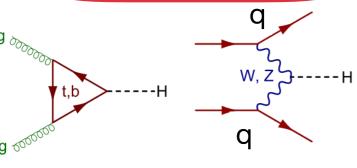
- Sufficient statistics for precision-level measurements.
- Path open to exploration of SM Lagrangian in the Eelectro-Weak symmetry braking sector.
- Probe to couplings to bosons and fermions
- Understand structure of its potential.



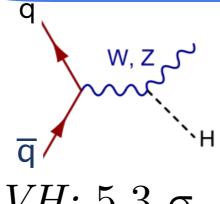




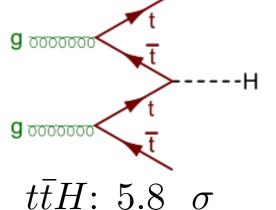








VH: 5.3 σ



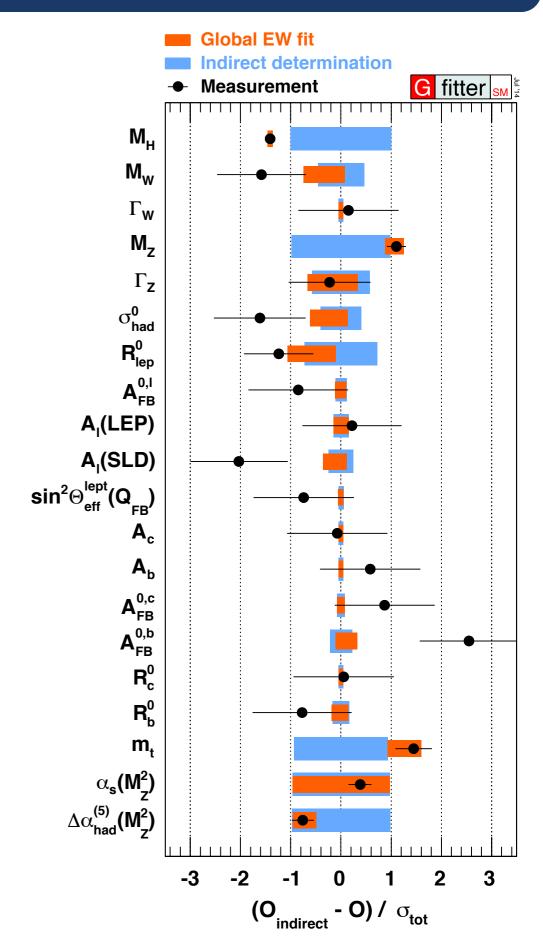
Introduction

Mass measurement

• The Higgs boson mass (m_H) is a fundamental free parameter of the Standard Model.

$$V(h) = \frac{1}{4}\lambda h^4 + \lambda vh^3 + \lambda v^2h^2$$

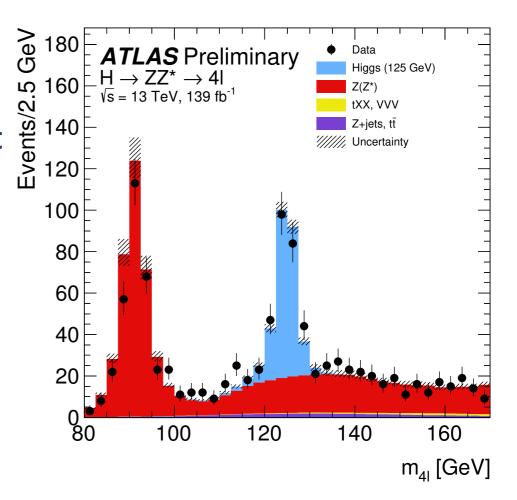
- Understanding the perturbative expansion of its potential $(\lambda v^2 h^2)$
- Precision determination allows for evermore precise higher order corrections to the cross section.
- Sensitivity to new physics in higher order corrections.
- Input to precision Electro Weak global fit.
- ▶ Key measurement of the LHC program.
- Aim in improving significantly on δm_H



Uncertainties

- For $\gamma\gamma$ and 4ℓ , signal is narrow resonant peak above a background continuum
 - ▶ Allows for precise Higgs boson mass measurement
 - ▶ Minimises the model dependency.
- Ingredients for optimal measurement of Higgs boson mass:
 - Detector performance driven measurement

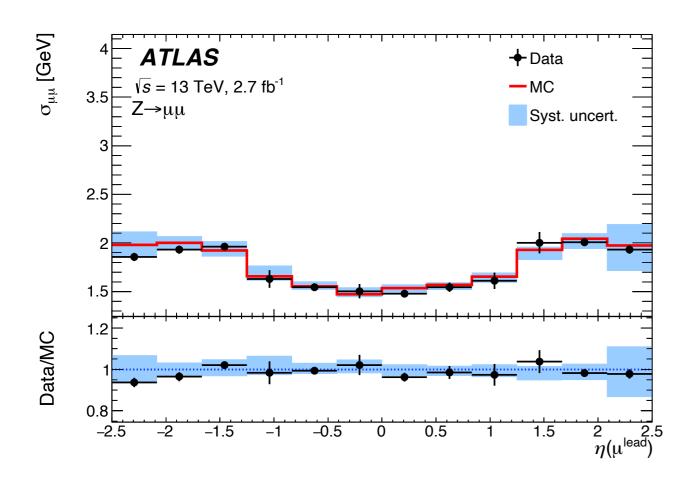
$$\delta m_H \simeq \frac{\sigma(m_{4\ell,\gamma\gamma})}{\sqrt{N-N_{\rm b}}}$$

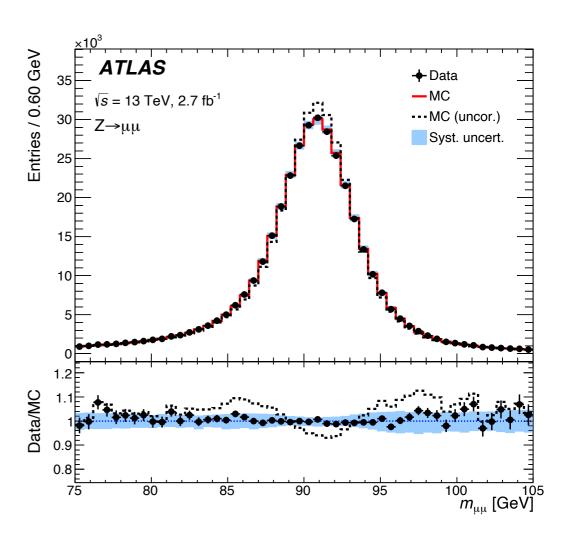


- (I) Statistical precision precision depends upon:
 - resolution of the reconstructed final state,
 - number of signal events.
- (II) Systematic uncertainty from understanding of detector performance:
 - energy and momentum scale,
 - resolution uncertainty.

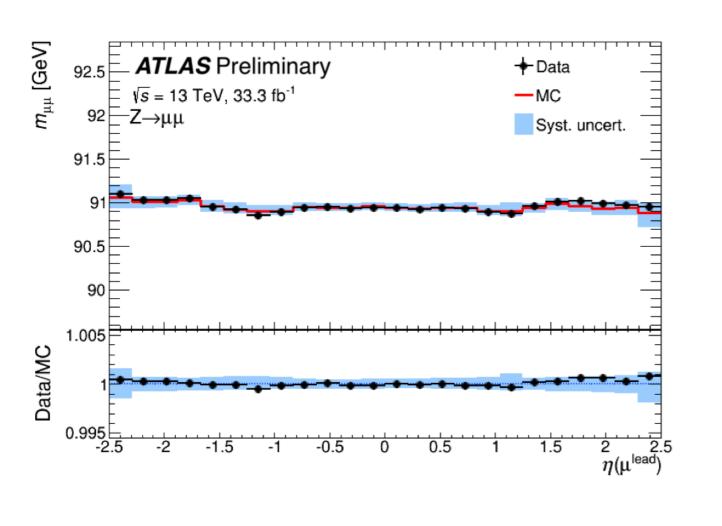


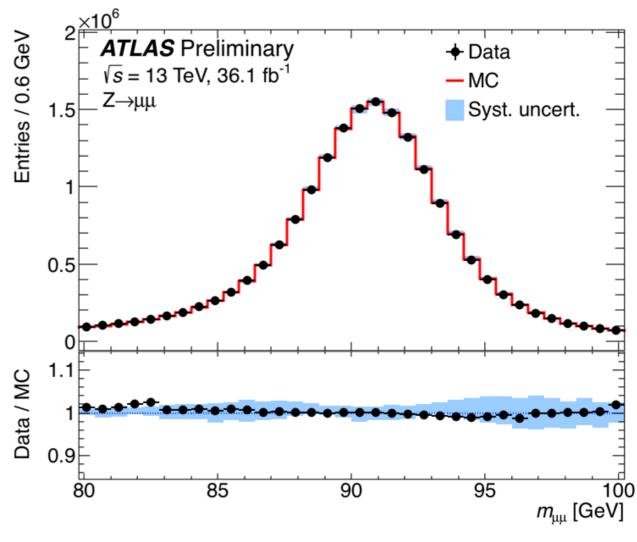
- Resolution muon channels (4μ , $2e2\mu$ and 4μ) crucial for m_H uncertainty:
 - ▶ Excellent momentum resolution of about 1% at about p_T 45 ~GeV.
- ullet Momenta calibrated to ${\mathrm J}/\psi$ and Z samples in data
 - for residual mis modelling of E^{loss} in calorimeters, alignment precision etc.
 - Including corrections to data accounting for alignment weak modes.
 - ▶ Precision down to 0.5 per mille for $|\eta|$ <1.0



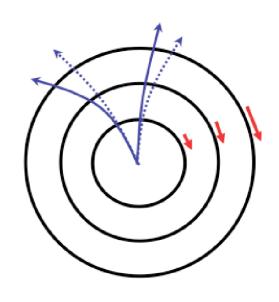


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- ullet Simulated momenta calibrated to ${
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 - for residual mis modelling of E^{loss} in calorimeters, alignment precision etc.
 - ▶ Uncertainty of about 10% on the resolution and 0.5% on the momentum scale.

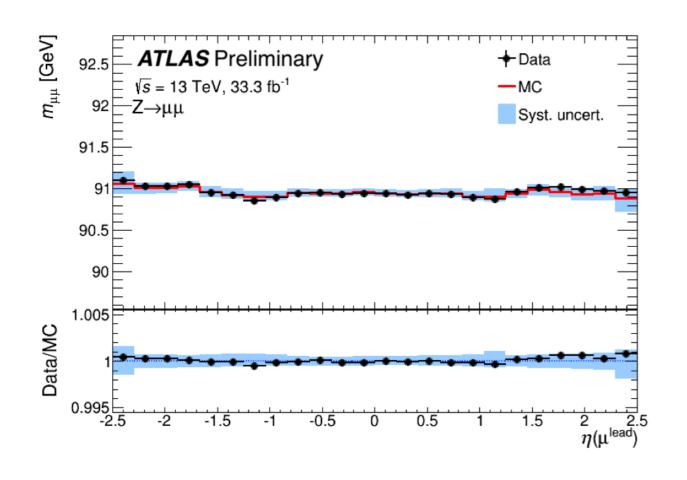


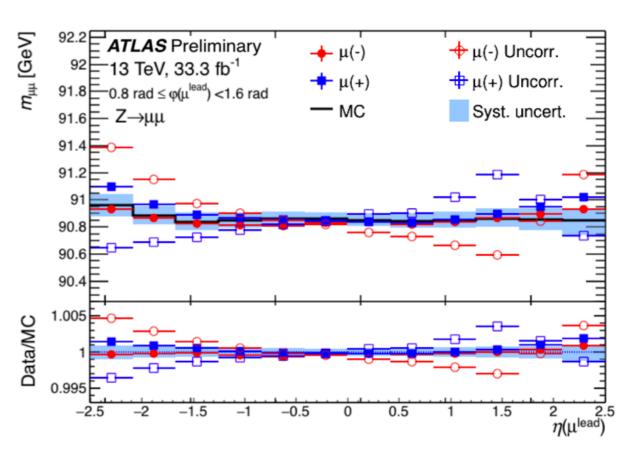


- Local misalignments and second order effects:
 - Charge dependent sagitta bias, with net effect of worsening resolution
 - In-situ correction based on $Z \rightarrow \mu\mu$ data, recovers up to 5% in resolution.



- Momentum scale understood down to the per mille level
 - ▶ Precision down to 0.5 per mille for $|\eta| < 1.0$

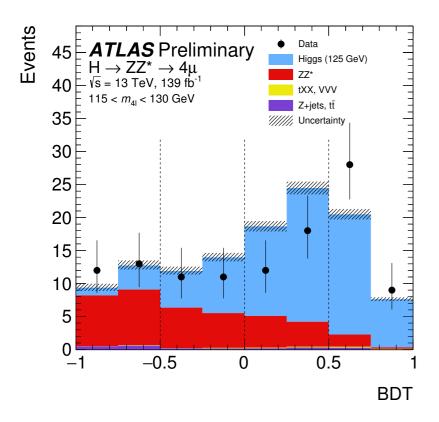


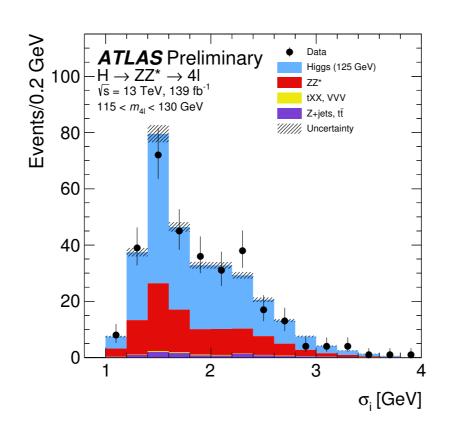


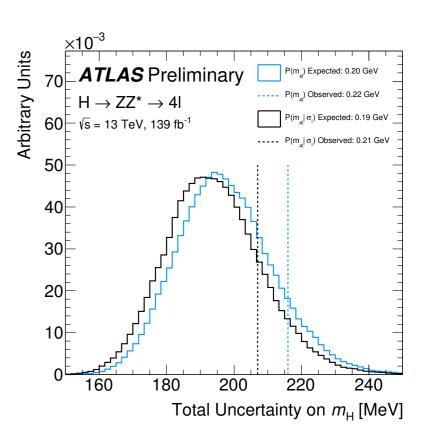
Three-prong approach to reduce uncertainty at analysis level:

arXiv:1806.00242

- (i) $\sim 15\%$ from m_{12} constraint to m_Z with kinematic fit and m_Z constraints on alignment weak modes.
- (i) ~2% from kinematic discriminant selecting signal and background events
 - ▶ Boosted Decision Tree on $p_T(4\ell)$, $y(4\ell)$ and $log(|\mathcal{M}_H|^2/|\mathcal{M}_{ZZ^*}|^2)$
- (ii) ~5% from multivariate per-event resolution likelihood.
 - Neural network to solve uncertainty correlations induced by kinematic discriminant.







Three-prong approach to reduce uncertainty at analysis level:

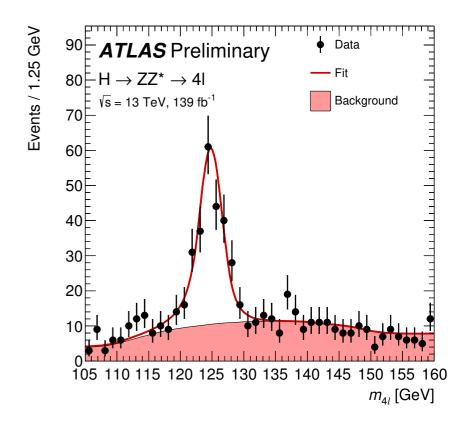
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Systematic uncertainty of ~70 MeV

Systematic Uncertainty	Impact (GeV)
Muon momentum scale	+0.08,-0.06
Electron energy scale	±0.02
Muon momentum resolution	±0.01
Muon sagitta bias correction	±0.01

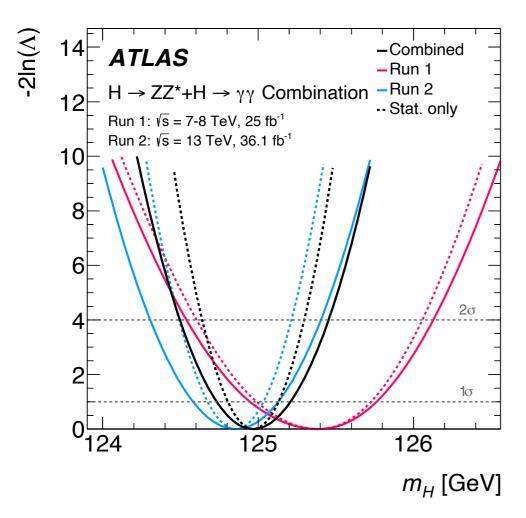
- ▶ 61% improvement w.r.t $m_{
 m H}{}^{H o ZZ,
 m Run \, I}$
- ▶ 15% improved precision w.r.t $m_{
 m H}^{
 m ATLAS+CMS,Run\,I}$

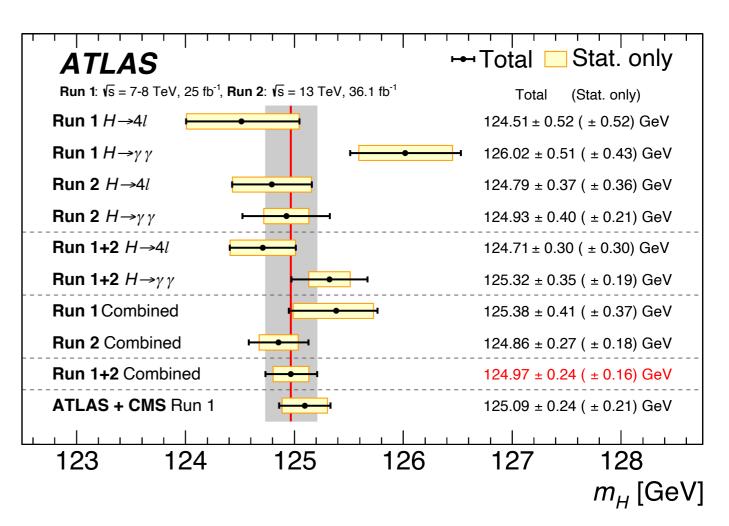


$$m_H = 124.92^{+0.21}_{-0.20} \text{ GeV}$$



• 4\ell and yy measurements are combined with ATLAS Run I result arXiv:1806.0024





Run 2 precision improved w.r.t Run 1.

$$m_H = 124.86 \pm 0.27(\pm 0.18 \text{ stat only}) \text{ GeV}$$

ATLAS Run I + 2 comparable precision to LHC Run I combination.

$$m_H = 124.97 \pm 0.24(\pm 0.16 \text{ stat only}) \text{ GeV}$$



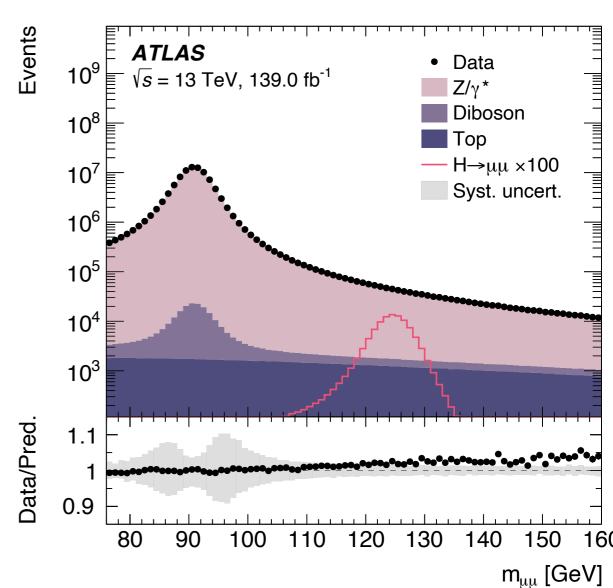
3.Decays to light leptons.

Second generation fermions

Light leptons

- Fermions acquire mass through Yukawa interactions with the Higgs field.
 - $+i\bar{\Psi}D\!\!\!/\psi$
 - Remains an elusive sector not probed by EWK precision tests.
 - $+D_{\mu}\Phi^{\dagger}D^{\mu}\Phi V(\Phi)$ $+(\bar{\Psi}_L\hat{Y}\Phi\Psi_R+h.c.)$

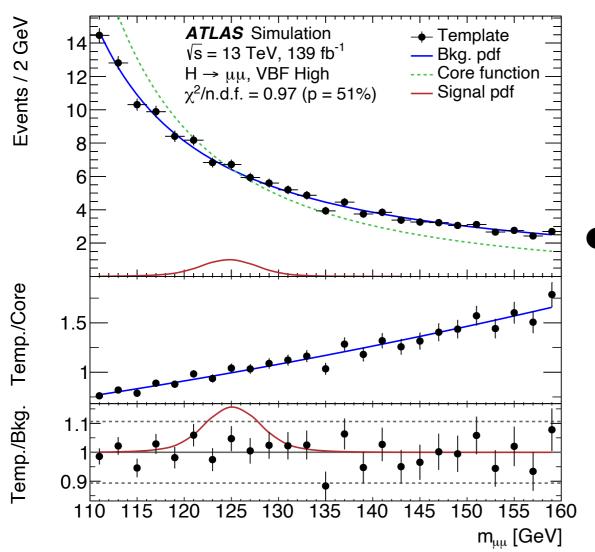
- ▶ LHC successfully probed couplings to third generation fermions ($ttH, H \rightarrow \tau \tau$)
- Next milestone, probe couplings to second and first generation fermions.
 - $H{\to}\mu\mu$ and ee offer unique insight.
 - ▶ Fully reconstructed final states with low hadronic activity.
 - Very rare processes:
 - → $\mathcal{B}(H \to \mu\mu)$ ~ (2.17 ± 0.04)×10⁻⁴
 - ◆ Large backgrounds from Drell-Yann production $Z \rightarrow \mu\mu$, ee

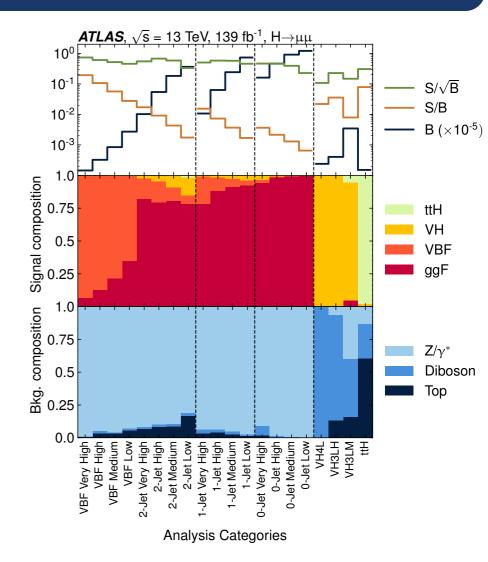


Second generation fermions

Light leptons

- Mutually exclusive categories
 - Targeting the various Higgs production modes, to increase s/b.
 - ► $H\rightarrow \mu\mu$, further splitting according process-specific multivariate boosted decision tree.
 - s/b ranging from 0.1% (0-jet Low) to 18% (VBF Very High)





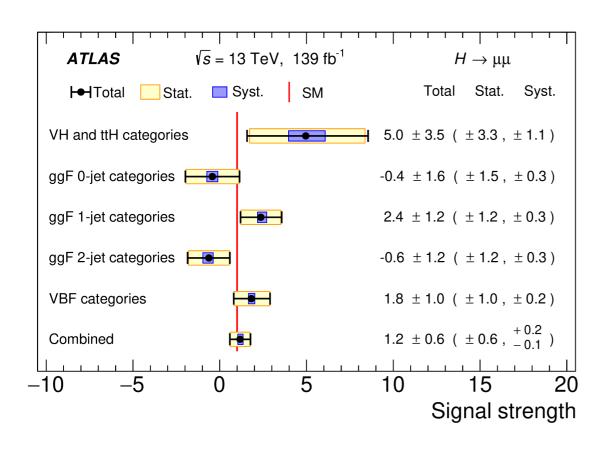
- Empirical background modelling in both analyses.
 - $F(m_{\mu\mu}) = \text{Rigid core}(m_{\mu\mu}) \times \text{Flexible Empirical}(m_{\mu\mu})$
 - Per-mille precision reached with two ad-hoc high statistics fast simulation ~10 ab-1 per category.

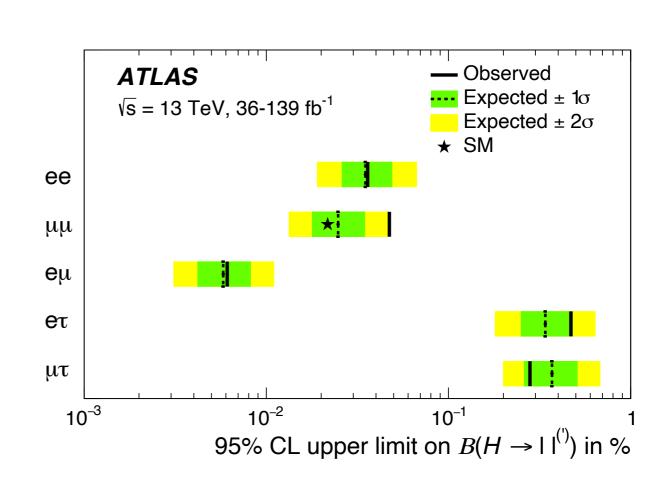
Results

Light leptons

 $\sqrt{s} = 13 \text{ TeV}, 139 \text{ fb}^{-1}$

- \bullet $\mathscr{B}(H{\to}ee) < 3.6 \times 10^{-4}$ at 95%CL.
 - \blacktriangleright Expected limit at $3.5 imes 10^{-4}$
 - ▶ Improvement of a factor of 5 w.r.t previous results.
- $H \rightarrow \mu\mu$: observed significance 2.0σ , $\mu = 1.2 \pm 0.5$.
 - Expected 1.7σ .
 - \bullet $\sigma(H \rightarrow \mu\mu) / \sigma^{SM}(H \rightarrow \mu\mu) < 2.2$ at 95%CL.
 - ▶ Improvement of a factor 2.5, with 25% from the methods used.





Entries / GeV



4. Production mode measurements

Production mode

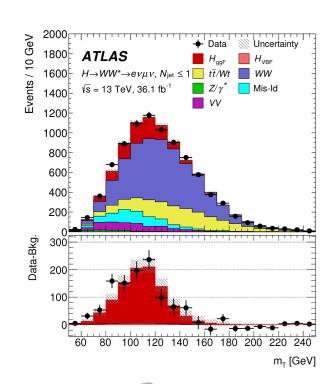
ATLAS full Run-2 combination

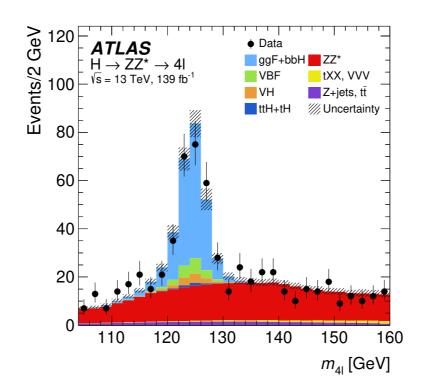
ATLAS-CONF-2020-027

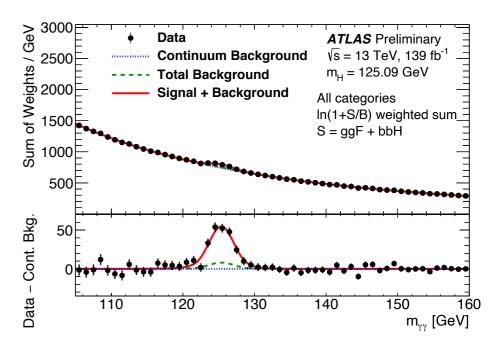
▶ Combined sensitivity of all channels to increase the precision of the Higgs

productions.

Analysis decay channel	Target Prod. Modes	\mathcal{L} [fb ⁻¹]	
$H o \gamma \gamma$	ggF, VBF, $WH, ZH, t\bar{t}H, tH$	139	
$H o ZZ^*$	ggF, VBF, WH , ZH , $ttH(4\ell)$	139	
$\Pi \rightarrow ZZ$	$t\bar{t}H$ excl. $H \rightarrow ZZ^* \rightarrow 4\ell$	36.1	
$H \to WW^*$	ggF, VBF	36.1	
$II \rightarrow VV VV$	$t ar{t} H$	30.1	
$H \to \tau \tau$	ggF, VBF	36.1	
$II \rightarrow 77$	t ar t H	30.1	
	VBF	24.5 - 30.6	
H o bar b	WH,ZH	139	
	$t ar{t} H$	36.1	
$H o \mu\mu$	${\rm ggF, VBF}, VH, t\bar{t}H$	139	
$H \rightarrow inv$	VBF	139	







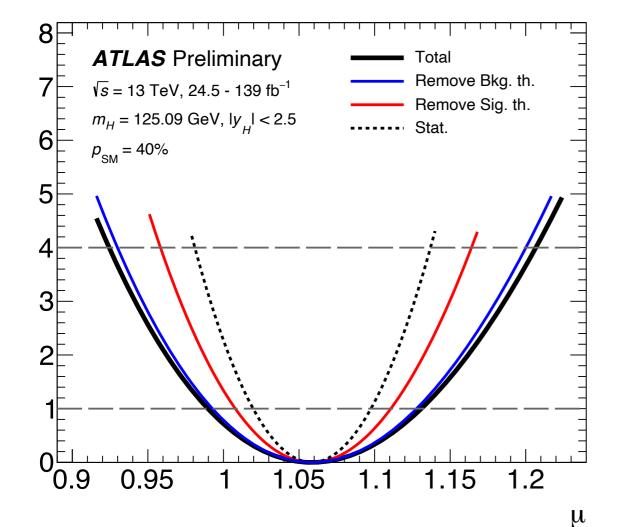
Simultaneous fit to all template cross sections

ATLAS-CONF-2019-005

- Extraction of global signal strength ($\mu = \sigma^{\text{obs}} / \sigma^{\text{exp}}$).
- ▶ Experimental sensitivity of the same oder as of theory (up to N³LO for ggF)

$$\mu = 1.06 \pm 0.07 = 1.06 \pm 0.04 \text{ (stat.)} \pm 0.03 \text{ (exp.)}^{+0.05}_{-0.04} \text{ (sig. th.)} \pm 0.02 \text{ (bkg. th.)}$$

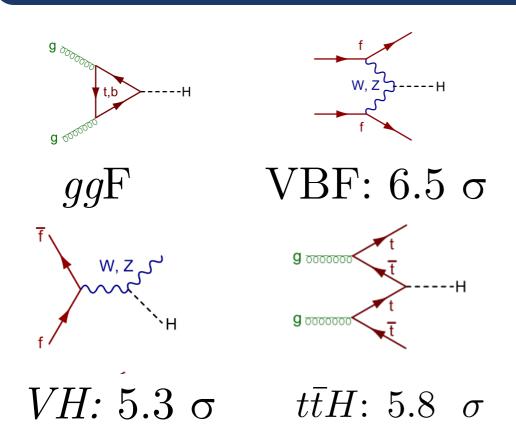
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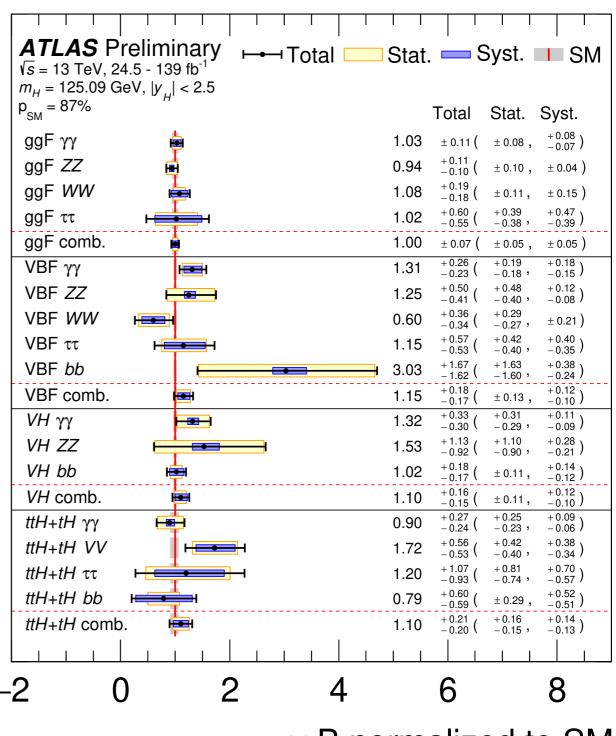
Uncertainty source	$\Delta\mu/\mu$ [%]
Luminosity	2.0
Background modeling	1.6
Jets, $E_{\mathrm{T}}^{\mathrm{miss}}$	1.4
Flavour tagging	1.1
Electrons, photons	2.2
Muons	0.2
au-lepton	0.4
Other	1.6
MC statistical uncertainty	1.7

Production mode measurements

Production modes

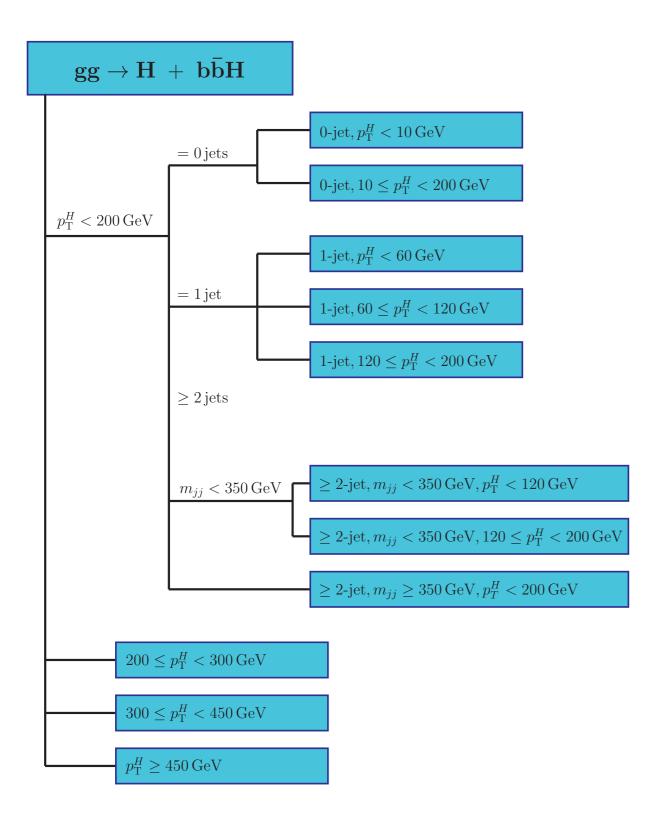


- In the first Run-2 data observed all SM production modes at the LHC.
 - ▶ With current precision uncertainties from 20% to 7% on production cross section.
 - Sufficient for more in-depth investigations into the couplings



 $\sigma \times B$ normalized to SM

Production modes

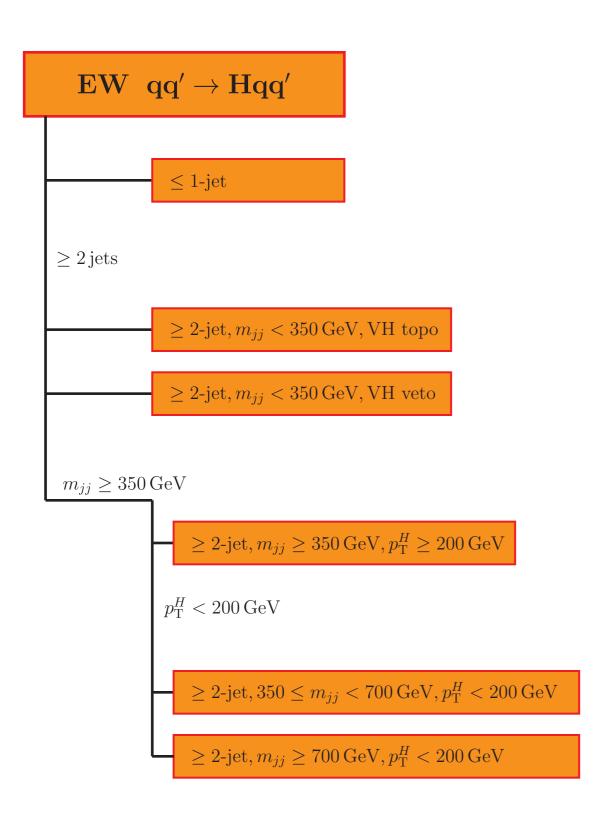


- Measure mutually exclusive phase spaces in agreement with theory and LHC experiments
 - In terms of kinematics of the Higgs or associated objects in productions.
 - Sensitivity to deviations from SM.
 - Avoidance of large modelling uncertainties.
 - Approximate experimental sensitivity.

- Advantage of complementary sensitivity in production from different final states:
 - $m_{\rm jj} > 450~{
 m GeV}$ from $H{
 ightarrow}WW^*$
 - $lackbox{High }p_{\mathrm{T}}{}^{H} ext{from }H{
 ightarrow}bb$



Production modes



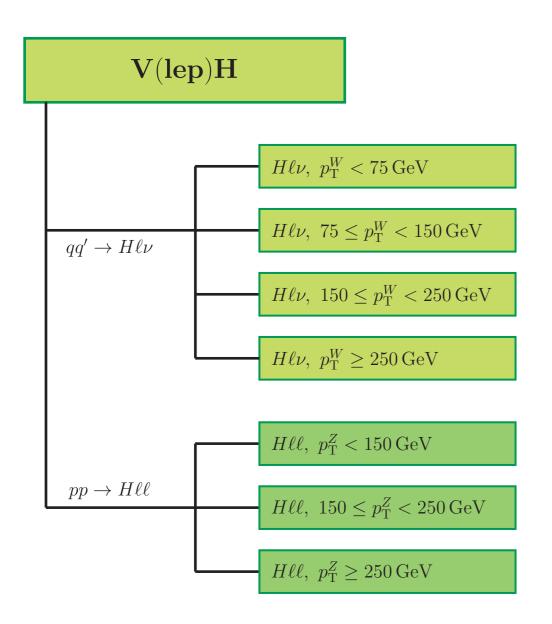
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Production modes

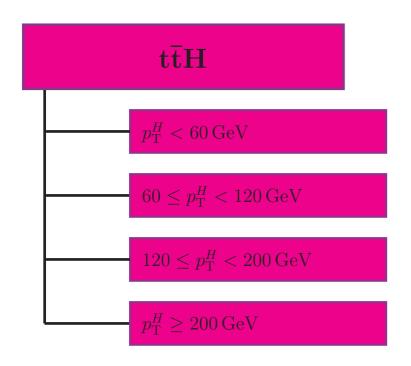


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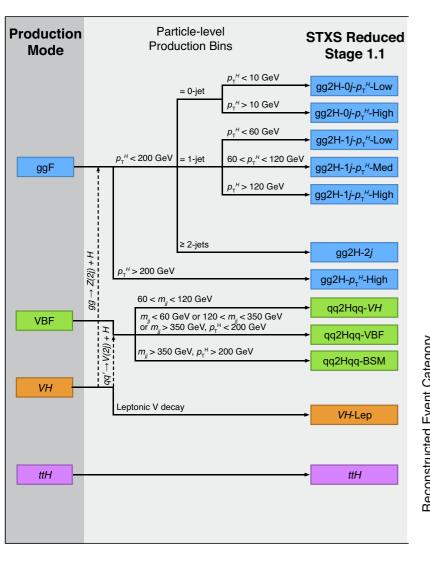


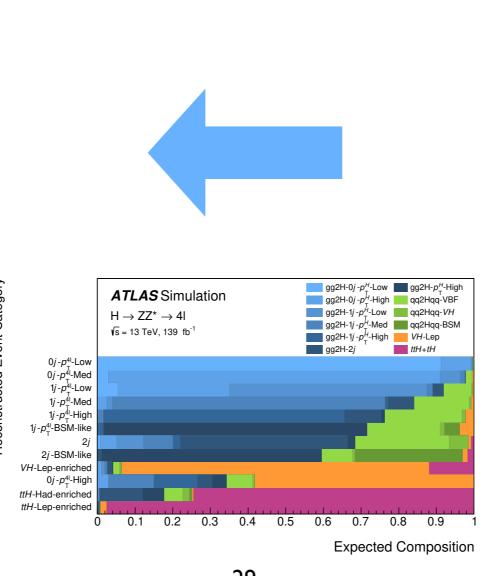
tH

Measurement strategy

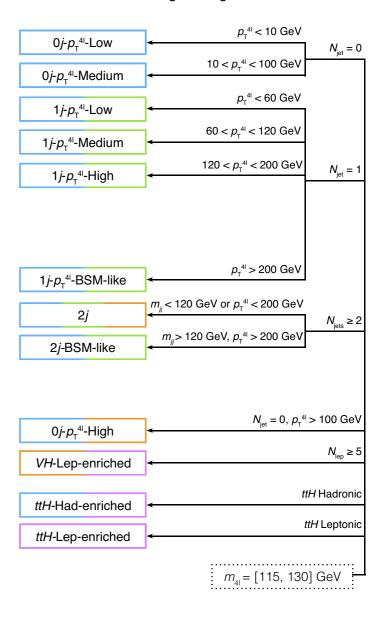
Production modes

- Strategy in measuring the cross section in these exclusive categories
 - lacktriangle Discussing here the example of the $H{ o}ZZ^*$
- Cut-based reconstruction-level categories,
 - maximising purity and minimising extrapolation to true phase-spaces.

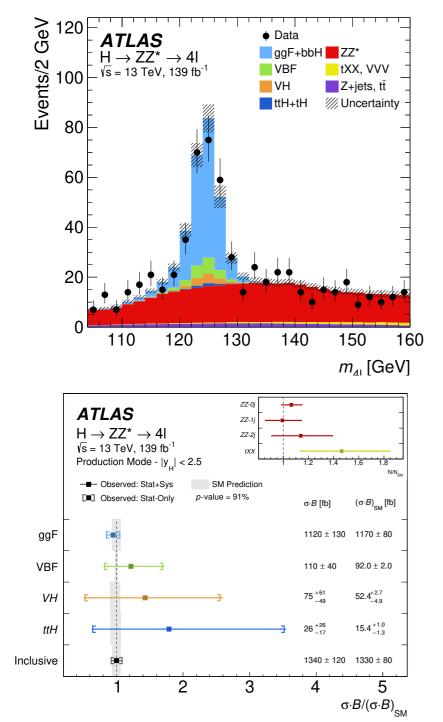


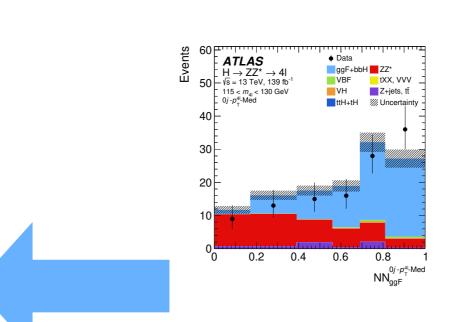


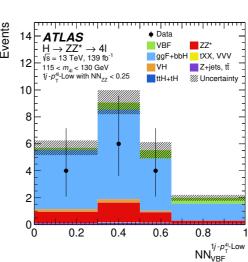
ATLAS √s = 13 TeV, 139 fb⁻¹ Reconstructed event categories Signal Region

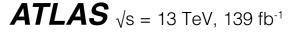


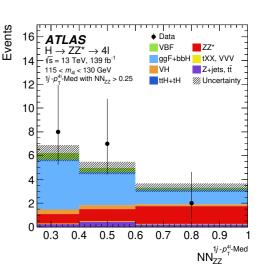
- Multi-output-node neural network discriminants in detector category.
 - \blacktriangleright Multidimensional fits on n-Idimensions on n output nodes per category.
- Backgrounds from data sidebands on resonant signal.
 - Performed as a function of the jet multiplicity to reduce higher order correction uncertainties.

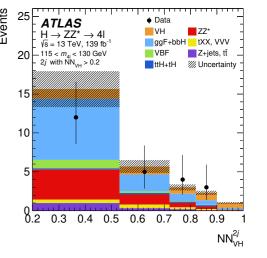










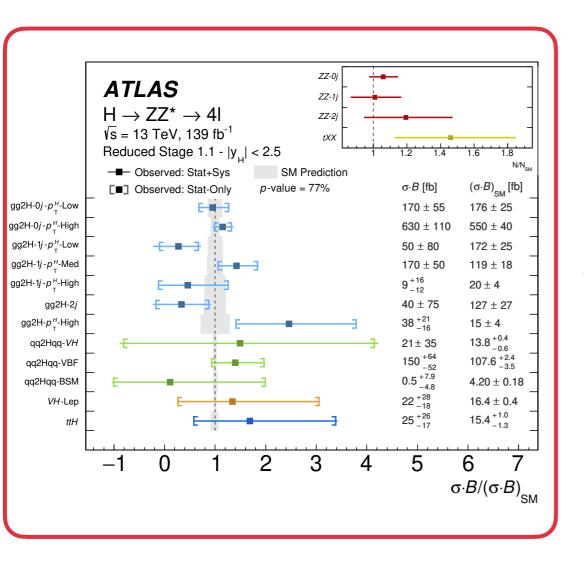


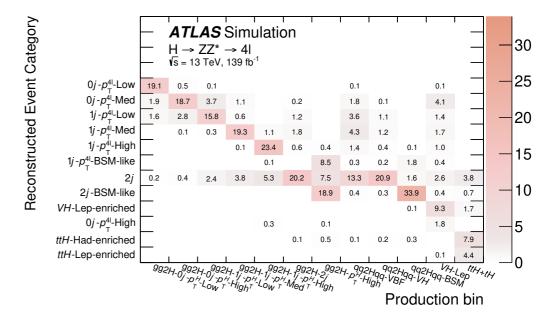
Production mode

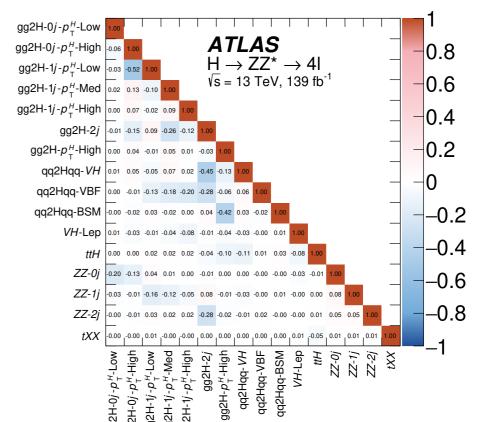
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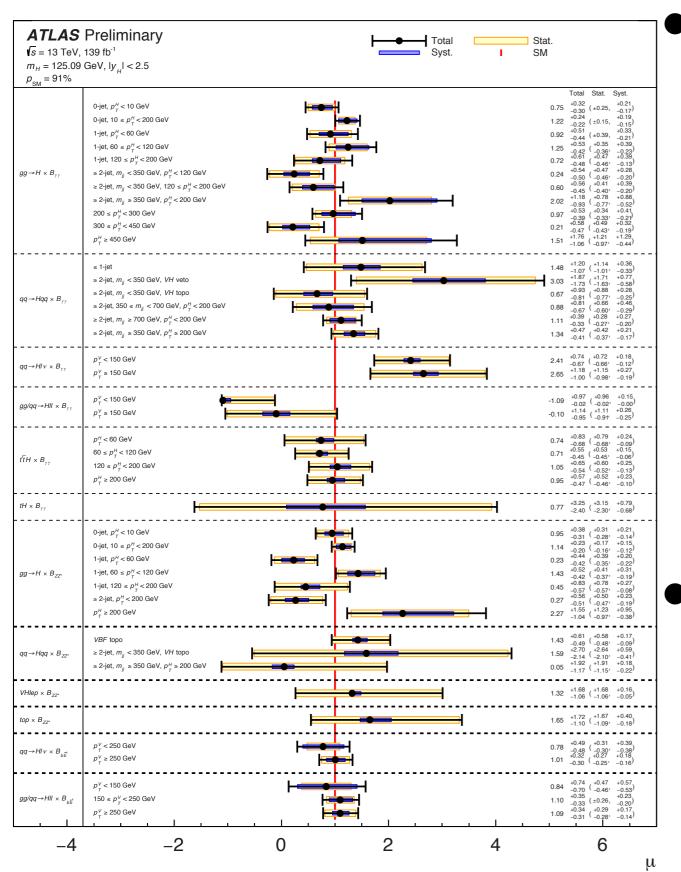








Production modes

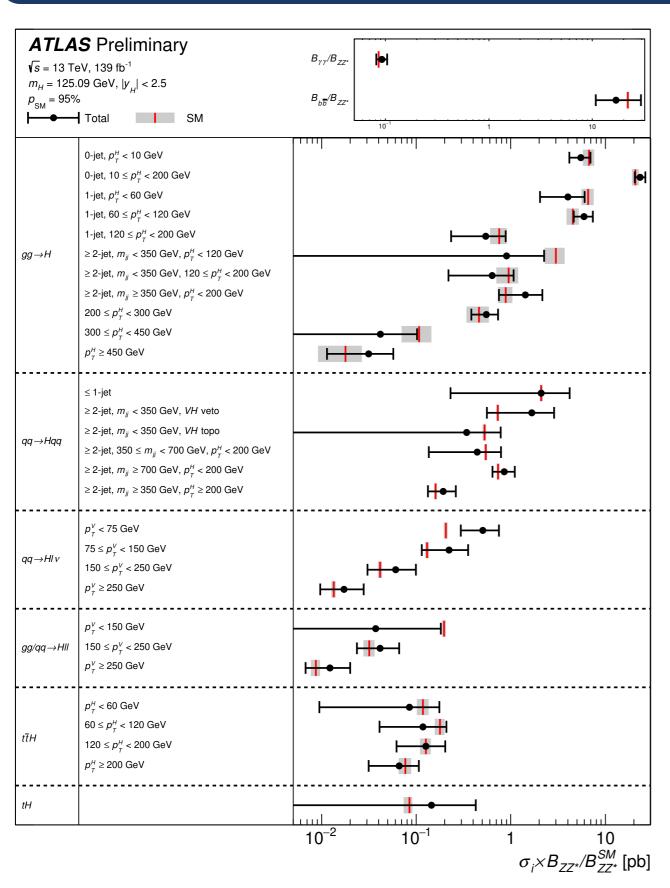


- Measure mutually exclusive phase-Spaces in agreement with theory and LHC experiments
 - In terms of kinematics of the Higgs or associated objects in productions.
 - Sensitivity to deviations from SM.
 - Avoidance of large modelling uncertainties.
 - Approximate experimental sensitivity.

- Advantage of complementary sensitivity in production from different final states:
 - $m_{\rm jj} > 450~{
 m GeV}$ from $H{
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 - $lackbox{High }p_{\mathrm{T}^{H}}\mathrm{from }H{
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Couplings interpretations

- Interpretation of couplings cross sections in the context of new physics.
- Assuming production and decay are factorised

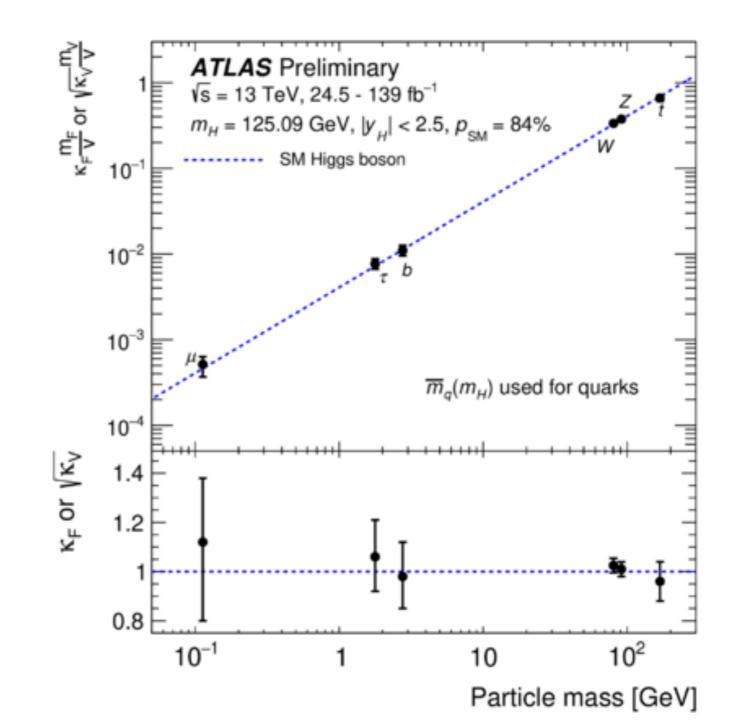
$$\sigma_i \times \mathbf{B}_f = \frac{\sigma_i(\kappa) \times \Gamma_f(\kappa)}{\Gamma_H}$$

Coupling strength modifiers

$$\kappa_j^2 = \frac{\sigma_j}{\sigma_j^{\text{SM}}} \qquad \kappa_j^2 = \frac{\Gamma_j}{\Gamma_j^{\text{SM}}}$$

$$\kappa_V = 1.03 \pm 0.03$$

$$\kappa_F = 0.97 \pm 0.07$$
.



ATLAS-CONF-2019-005

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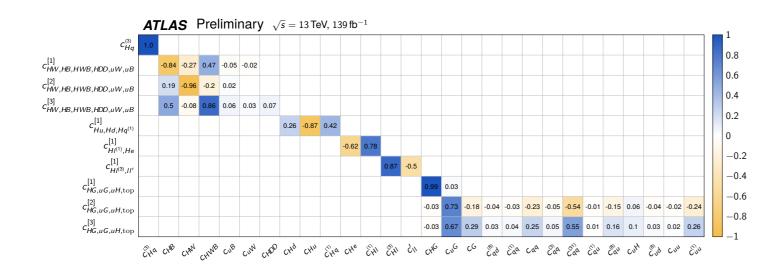
Couplings interpretations

Production mode

Results interpreted in the context of new physics:

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{i}^{N_{d6}} \frac{c_i}{\Lambda^2} O_i^{(6)} + \sum_{j}^{N_{d8}} \frac{b_j}{\Lambda^4} O_j^{(8)} + \dots$$

- Standard Model Effective Field Theory as the standard candle.
- Probe for non-SM contributions to the tensor structure of the Higgs boson.
- Enhance sensitivity
 - by isolating dependencies in Wilson coefficients (c_i) allowing for simultaneous extraction through eigenvector decomposition of the dependencies.



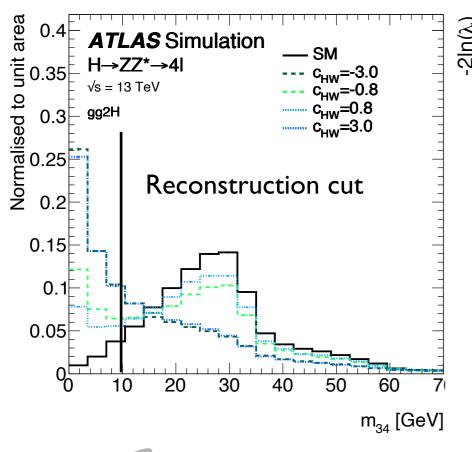
Coefficient	Operator	Example process
c_{HDD}	$\left(H^\dagger D^\mu H\right)^* \left(H^\dagger D_\mu H\right)$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$
c_{HG}	$H^\dagger H G^A_{\mu u} G^{A\mu u}$	g g
c_{HB}	$H^\dagger H B_{\mu u} B^{\mu u}$	$ \begin{array}{ccccc} & q & & q & \\ & Z & & & H & \\ & q & & Z & & q \end{array} $
c_{HW}	$H^\dagger H W^I_{\mu u} W^{I \mu u}$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$
c_{HWB}	$H^\dagger au^I H W^I_{\mu u} B^{\mu u}$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$
c_{eH}	$(H^{\dagger}H)(\bar{l}_{p}e_{r}H)$	$H \longrightarrow \stackrel{\ell}{\longleftarrow}_{\ell}$
$c_{Hl}^{ ext{ iny (1)}}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{l}_{p}\gamma^{\mu}l_{r})$	$ \begin{array}{c} q \\ q \end{array} $ $ \begin{array}{c} Z \\ \ell \\ H \end{array} $
$c_{Hl}^{\scriptscriptstyle (3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{l}_{p}\tau^{I}\gamma^{\mu}l_{r})$	$q \longrightarrow W \stackrel{\nu}{\underset{\ell}{\swarrow}} \stackrel{\nu}{\underset{H}{\swarrow}}$
c_{He}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{e}_{p}\gamma^{\mu}e_{r})$	q Z e e H
$c_{Hq}^{{\scriptscriptstyle (1)}}$	$(H^\dagger i \overleftrightarrow{D}_\mu H) (\bar{q}_p \gamma^\mu q_r)$	q Z ℓ ℓ H
$c_{Hq}^{\scriptscriptstyle (3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{q}_{p}\tau^{I}\gamma^{\mu}q_{r})$	q V
c_{Hu}	$(H^\dagger i \overleftrightarrow{D}_\mu H) (\bar{u}_p \gamma^\mu u_r)$	$u \xrightarrow{Z} \ell \\ u \xrightarrow{H}$
c_{Hd}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{d}_{p}\gamma^{\mu}d_{r})$	$\begin{array}{c c} d & Z & \ell \\ d & & H \end{array}$

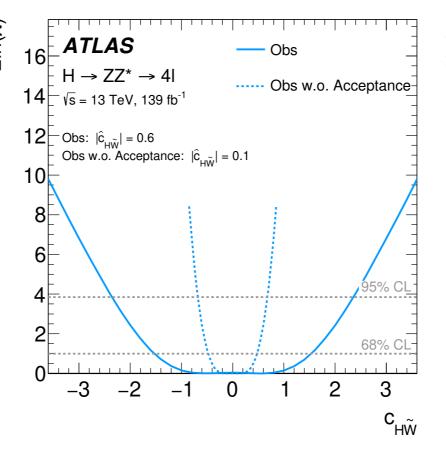


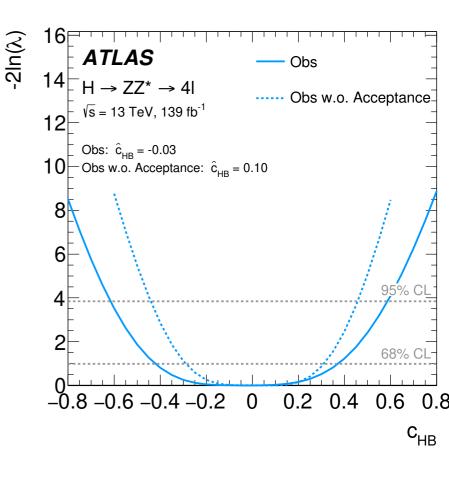
Results interpreted in the context of new physics:

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- ▶ Standard Model Effective Field Theory as the standard candle.
- ▶ Probe for non-SM contributions to the tensor structure of the Higgs boson.
- Account for BSM acceptance effects in kinematic observables of decay products



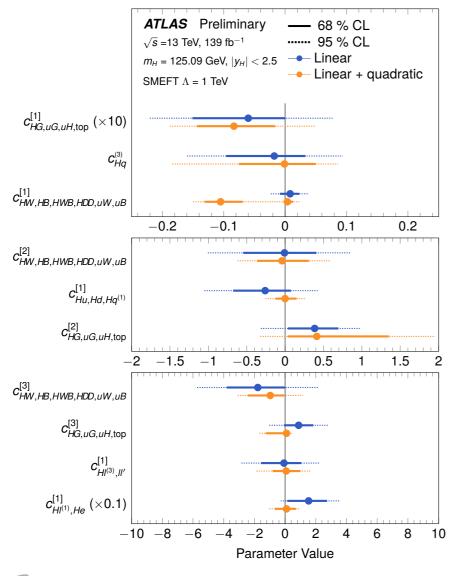


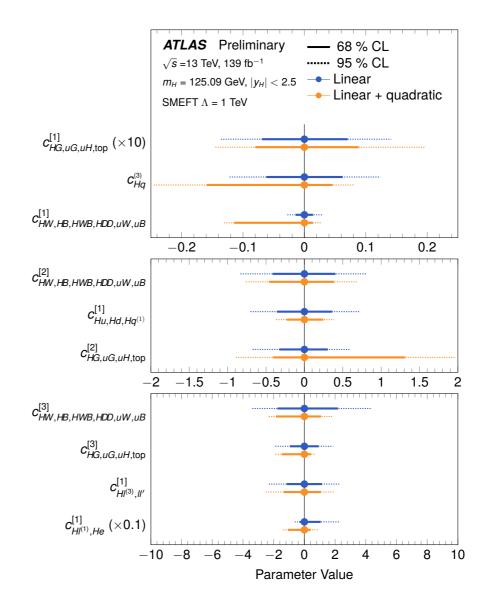


- Results interpreted in the context of new physics:
 - ▶ Results in both

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{i}^{N_{d6}} \frac{c_i}{\Lambda^2} O_i^{(6)} + \sum_{j}^{N_{d8}} \frac{b_j}{\Lambda^4} O_j^{(8)} + \dots$$

- ◆ linear approximation for dim-6 operators and,
- linear plus quadratic approximation for general sensitivity to dim-8, suppressed by Λ-4





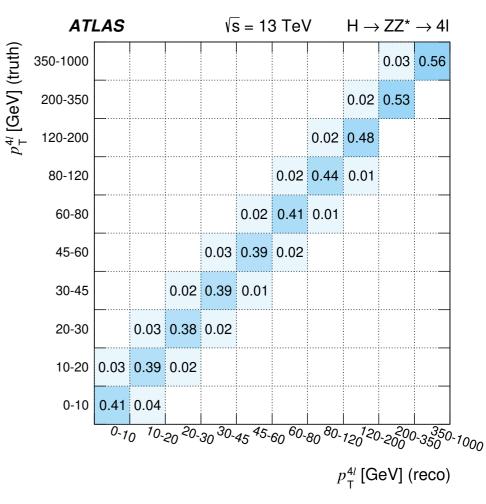
5. Differential cross section measurements

Differential cross section

- At Run II sufficient statistics for constraining differential measurements
- Fiducial cross section definition
 - lacktriangledown including detector efficiency (C), detector acceptance (A) and branching ${\mathscr B}$

$$\sigma_{i,\text{fid}} = \sigma_i \times A_i \times \mathcal{B} = \frac{N_{i,\text{fit}}}{\mathcal{L} \times C_i}$$

- ▶ Cuts mimicking reconstruction selection:
 - (i) Model independent result.
 - (ii) No extrapolation beyond measurable phase-space
- In diboson channels, resonant peak over smooth background
 - ▶ Good resolution on final-state particles, in particular in $H{\to}4{\it \ell},\ \gamma\gamma$
- Unfolding performed within the signal extraction fit
 - ▶ Via detector response inversion.
 - Reduces further any model assumptions in disentangling for detector effects



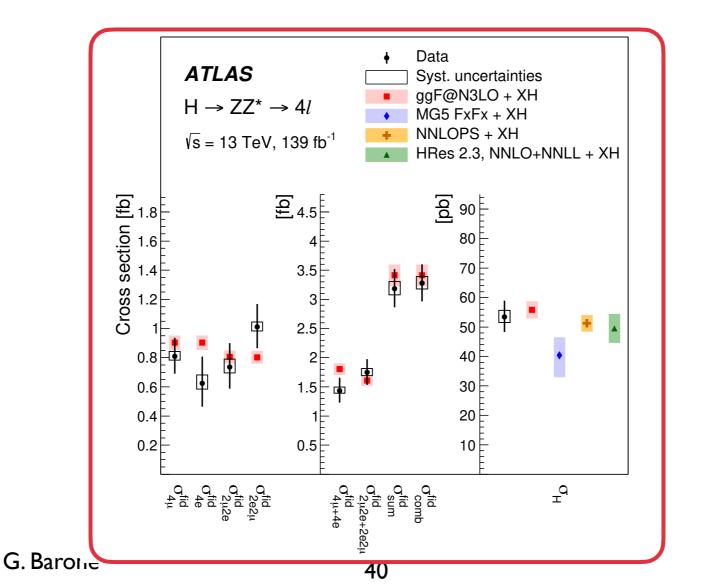
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Differential cross section

• Inclusive fiducial cross sections:

$$H\rightarrow\gamma\gamma$$
 60.4 ± 6.1 (stat.) ± 6.0 (exp.) ± 0.3 (theo.) fb

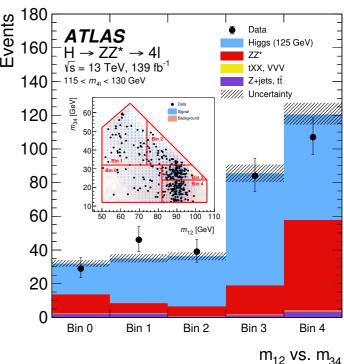
- \blacktriangleright SM predictions of 3.33 \pm 0.15 fb ($H{\to}ZZ$) and 63.5 \pm 3.3 fb ($H{\to}\gamma\gamma$)
- ullet For ZZ also cross section per final state
 - ▶ Eventually sensitivity to final state interference (10%) in same flavour quadruplets

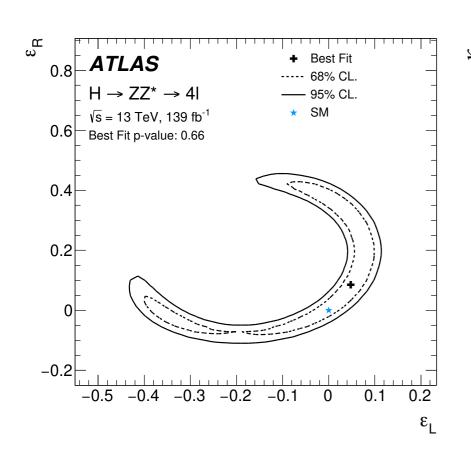


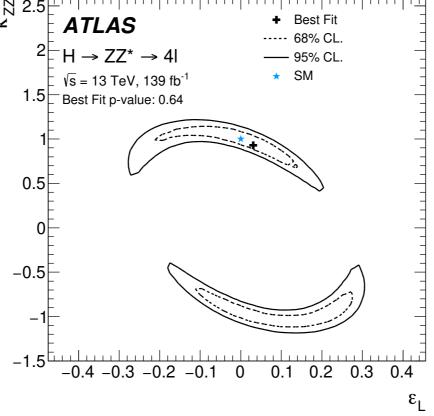


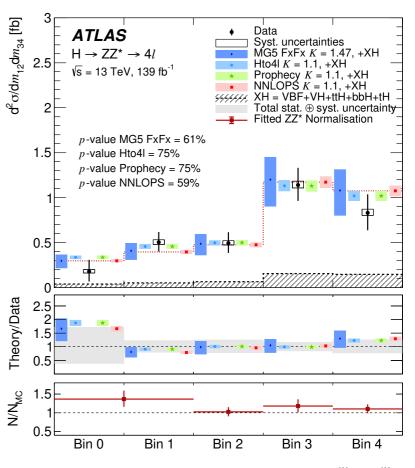
Interpretation

- In $H \rightarrow 4\ell$ m_{12} vs m_{34} : sensitivity to contact interactions:
 - \triangleright ε_{R} , ε_{L} and K: flavour universal modifiers of the contact terms between H, Z and leptons (arXiv:1504.04018)
 - ◆ Angular distributions unaffected: same Lorentz structure as SM term.





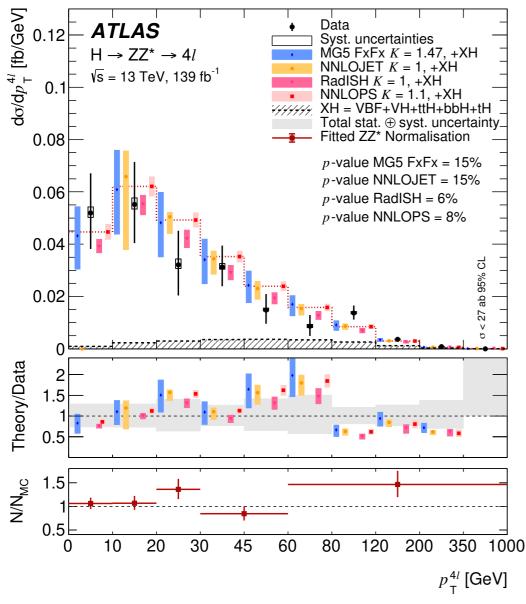


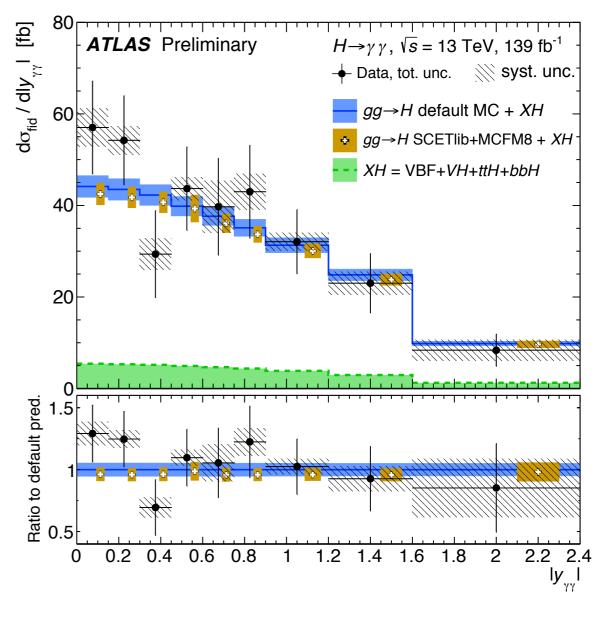


 $m_{12} \text{vs.} m_{34}$

Higgs boson kinematics

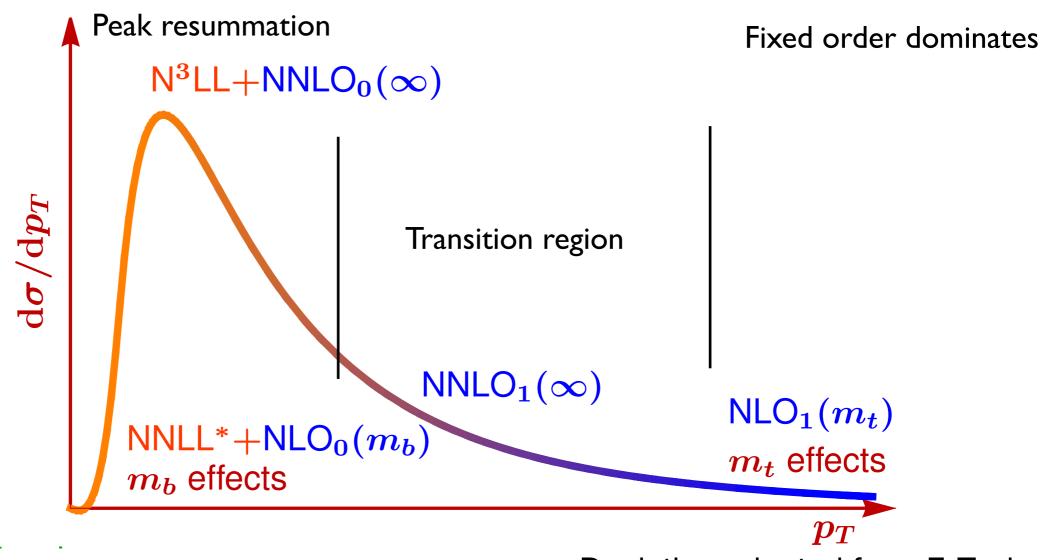
- 4ℓ : isolate signal under the Higgs resonant peak (115 < $m_{4\ell}$ < 130).
- ullet $\gamma\gamma$ cross section extracted from resonant peak over the $\gamma\gamma$ continuum.
- Higgs boson $p_{T,4\ell(\gamma\gamma)}$ and rapidity $(y_{4\ell(\gamma\gamma)})$ probe.
 - $p_{T,4\ell(\gamma\gamma)}$: Lagrangian structure of H interactions, Yukawa couplings
 - $y_{4\ell(\gamma\gamma)}$: Sensitivity to proton's parton density functions.





Differential cross section

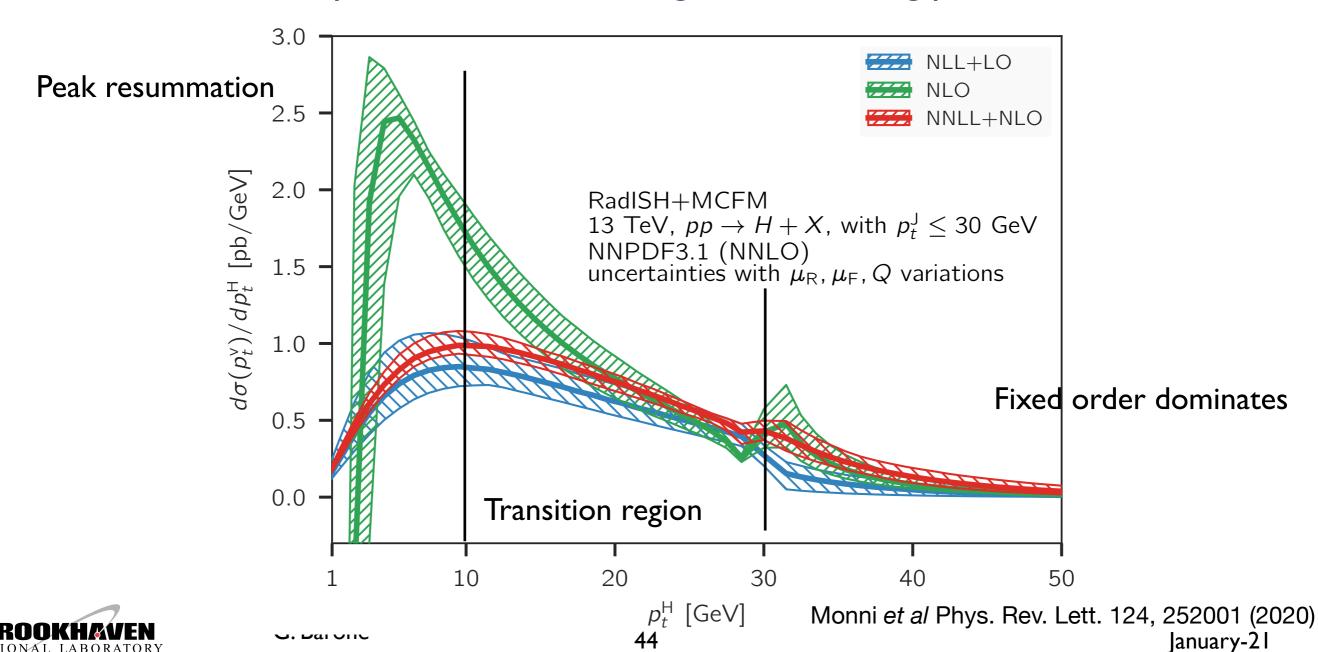
- At Run II sufficient statistics for constraining differential measurements
 - Increased precision needed to disentangle effects from higher-order corrections from observables spectra:
 - \blacktriangleright Ex Higgs p_T as a function of jet vetos





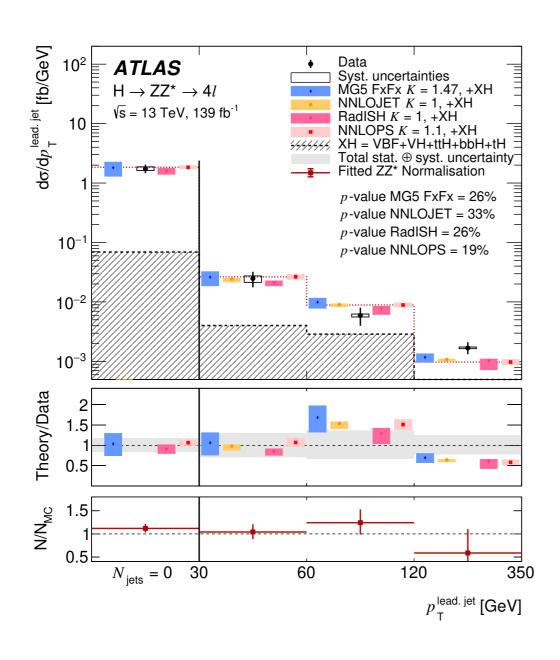
Depiction adapted from <u>F. Tackmann</u>

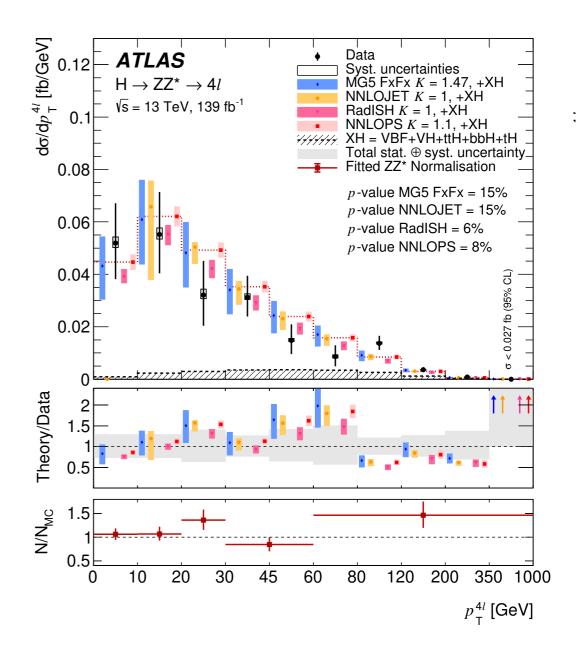
- At Run II sufficient statistics for constraining differential measurements
 - Increased precision needed to disentangle effects from higher-order corrections from observables spectra:
 - \blacktriangleright Ex Higgs p_T as a function of jet vetos
 - ▶ State of the art predictions in these regions start being published.



H and jets variables

- Increased precision needed to disentangle effects from higher-order corrections from observables spectra:
 - ▶ Measurements at Run-2 competitive on these state-of-the art predictions
 - ▶ Ex RADISH and NNLOJET.

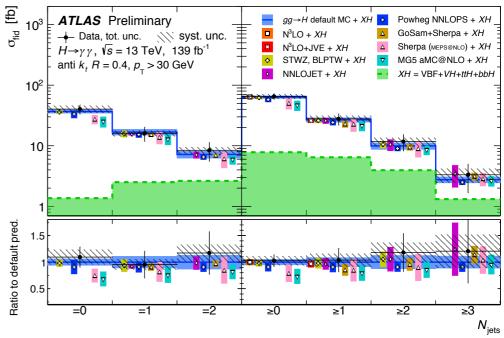


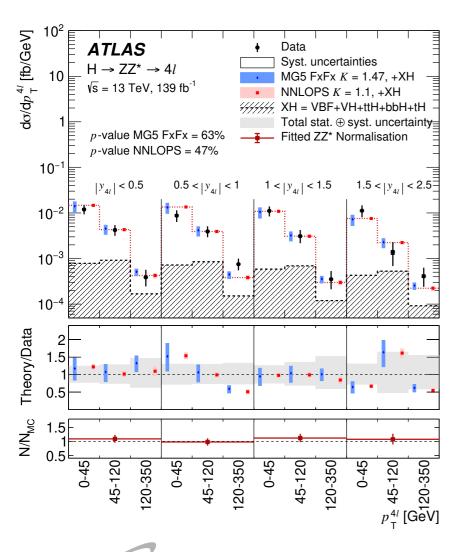


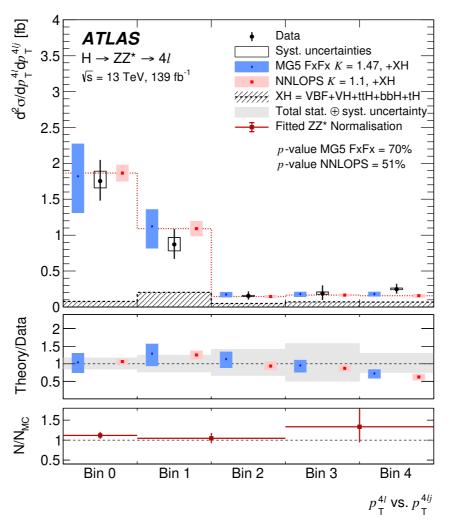


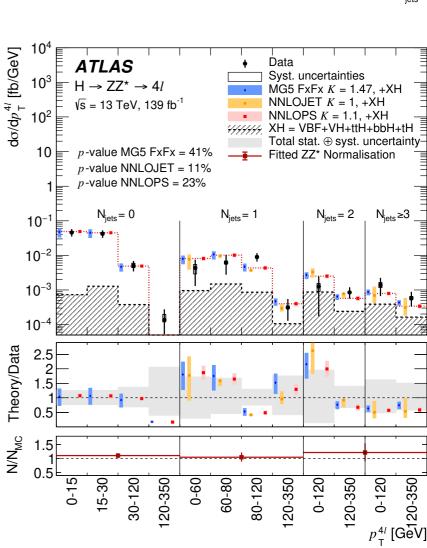
H and jets variables

- Double differential $d^2\sigma/dp_{\rm T}dN_{\rm J}$:
 - Probe the Higgs production mode
 - ▶ $N_j = 0$ dominated buy ggF production
 - ▶ Nj > I VBF enriched production





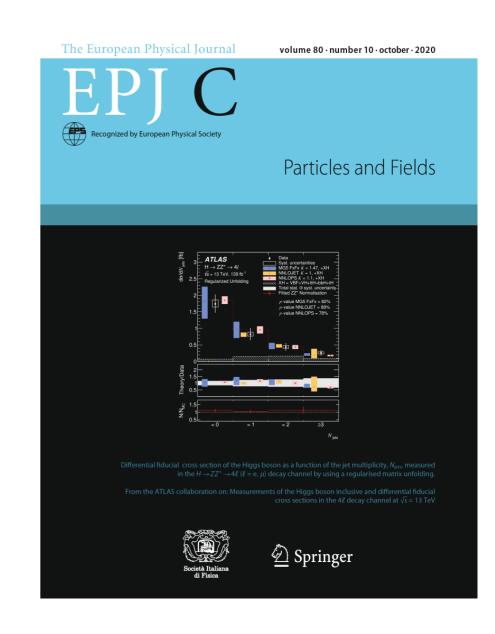




6. Conclusions

Conclusion

- Higgs physics provide an excellent picture for
 - ▶ Searches for new phenomena resonant at higher scales.
 - ▶ Searches for deviations to theory within the scales of the experiment.
- Run II first results of ATLAS in the study of Higgs boson properties
 - 1. Measurement of m_H at 2 per mille precision level.
 - 2. Fiducial cross section measurements, sensitivity to several distributions
 - 3. Production mode analysis and template cross section measurements.



Presented only a selection results full set

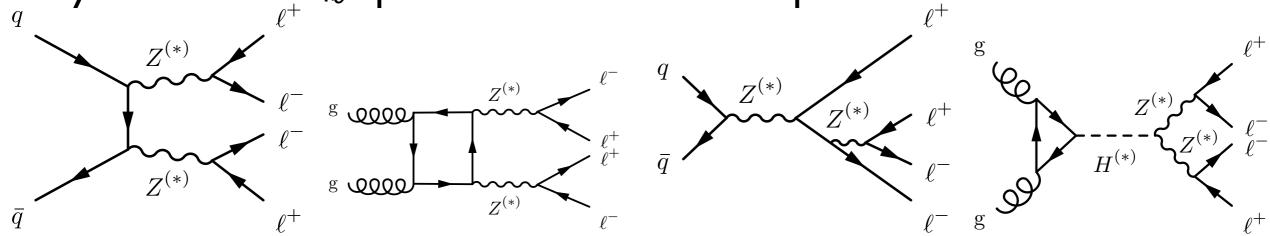
https://twiki.cern.ch/twiki/bin/view/AtlasPublic/HiggsPublicResults



Additional material

4ℓ mass spectrum

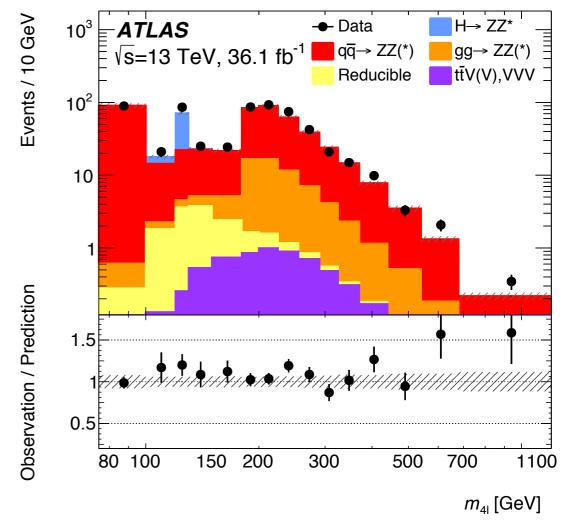
• Study of the the $m_{4\ell}$ spectrum and Offshell H production



- $m_{4\ell}$ ranges from single Z resonance, including H production up to ZZ production
- ▶ Extraction of the BR($Z \rightarrow 4\ell$)

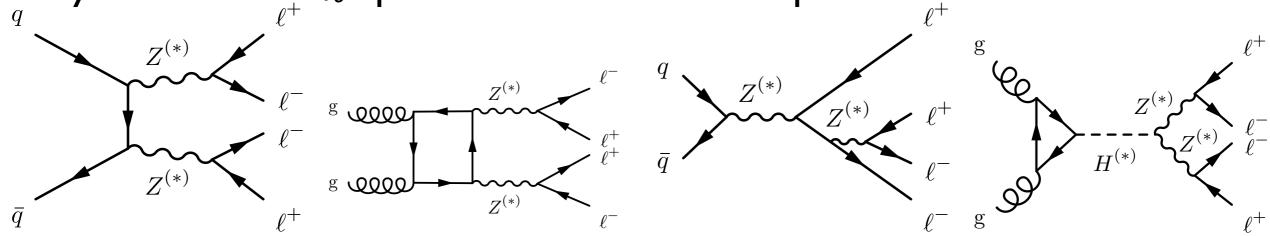
Measurement	$\mathcal{B}_{Z\to4\ell}/10^{-6}$
ATLAS, $\sqrt{s} = 7$ TeV and 8 TeV [8]	$4.31\pm0.34(stat)\pm0.17(syst)$
CMS, $\sqrt{s} = 13 \text{ TeV } [6]$	$4.83^{+0.23}_{-0.22}(\text{stat})^{+0.32}_{-0.29}(\text{syst})\pm0.08(\text{theo})\pm0.12(\text{lumi})$
ATLAS, $\sqrt{s} = 13 \text{ TeV}$	$4.70 \pm 0.32(\text{stat}) \pm 0.21(\text{syst}) \pm 0.14(\text{lumi})$

- Offshell Higgs production, enhanced at 350
 GeV because of top-quark loops in ggF
 - ◆ Including interference between H and ZZ productions.
- Above ~ $2m_Z$ enhancements of $qq \rightarrow ZZ$ and $gg \rightarrow ZZ$.



40 mass spectrum

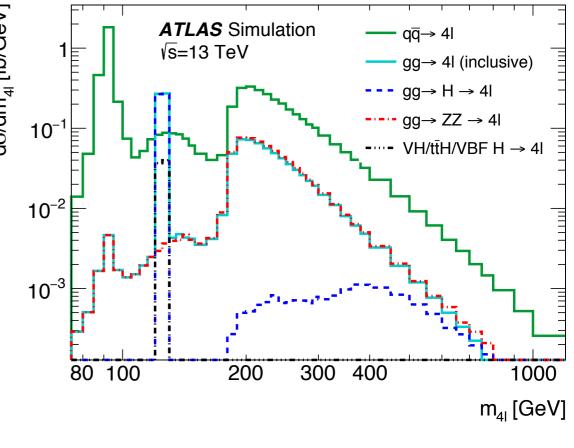
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ATLAS, $\sqrt{s} = 7$ TeV and 8 TeV [8]	$4.31\pm0.34(stat)\pm0.17(syst)$	Ū,
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		- て

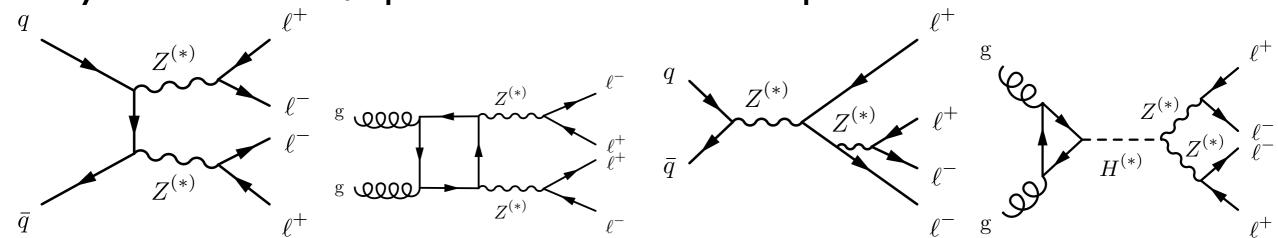
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January-21

4ℓ mass spectrum

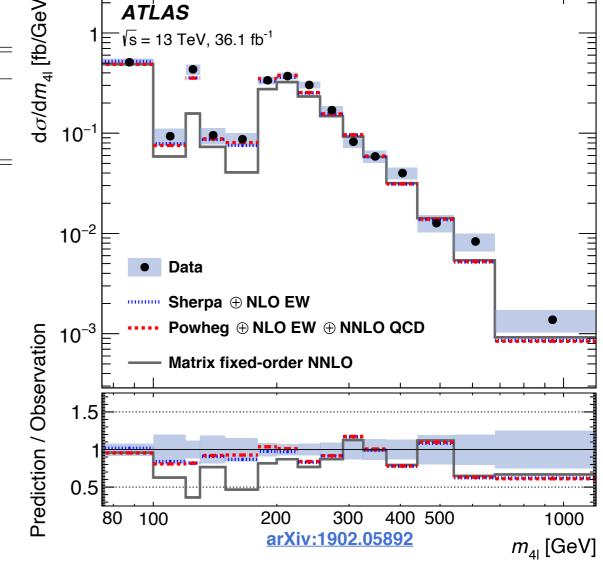
• Study of the the $m_{4\ell}$ spectrum and Offshell H production



- \blacktriangleright $m_{4\ell}$ ranges from single Z resonance, including H production up to ZZ production
- ▶ Extraction of the BR($Z \rightarrow 4\ell$)

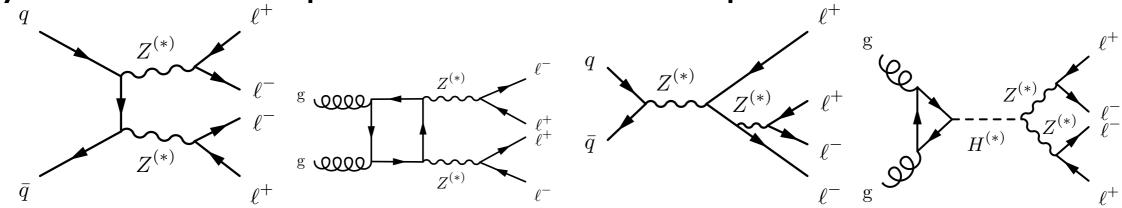
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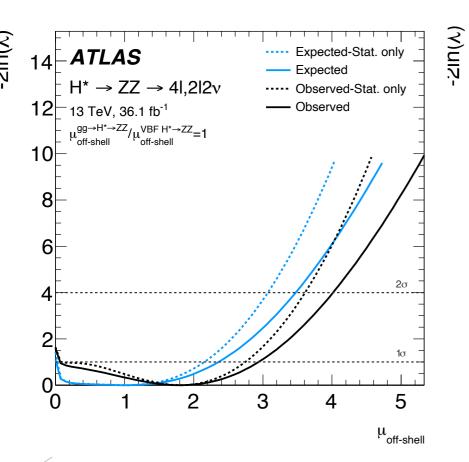


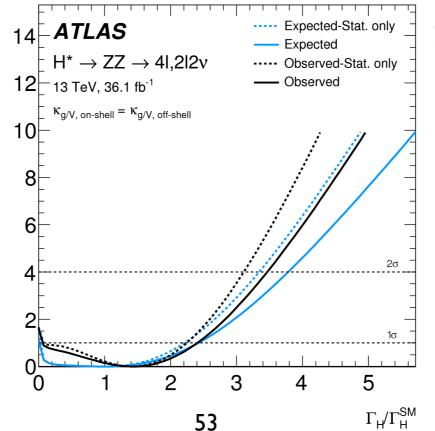
• Study of the the $m_{4\ell}$ spectrum and off-shell H production

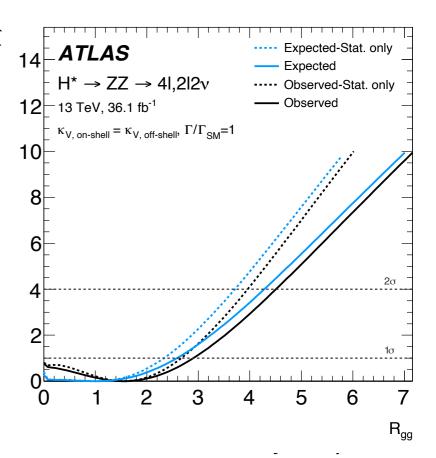
arXiv:1902.05892



- Offshell Higgs production, enhanced at 350 GeV because of top-quark loops in ggF
- Measured upper limit on width combining 4ℓ and $\ell \overline{\ell} \nu \overline{\nu}$
- Limit Γ_H possibile from the off-shell to on-shell event yield ratio R_{gg}
 - ullet on-shell event yields $\sim k^2 {
 m g,on-shell} \ / \ \Gamma_H$, while off-shell $\sim k^2 {
 m g,off-shell}$





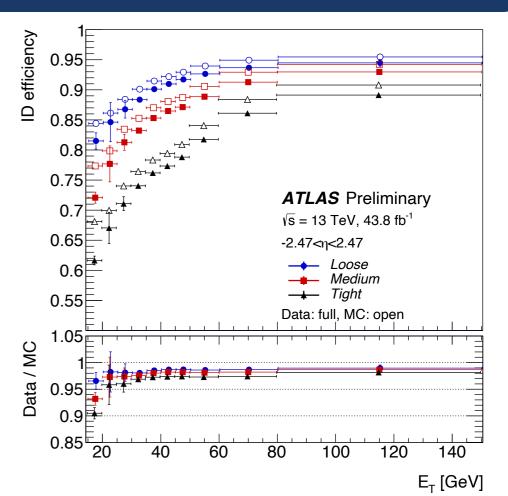


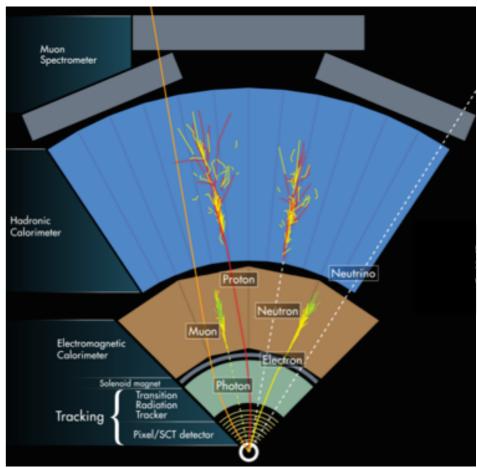
Reconstruction and selection

- Electrons (e).
 - Isolated objects clustered from calorimeter energy deposits with associated ID track.
 - ▶ $E_T > 7$ GeV, $|\eta| < 2.47$ and $|z_0 \sin(\vartheta)| < 0.5$ mm
- Muons (μ) .
 - Combined track fit of Inner Detector and Muon Spectrometer hits,
 - ▶ $p_T > 5$ GeV, $|\eta| < 2.7 |z_0 \sin(\vartheta)| < 0.5$ mm of "loose or medium quality"
 - Isolated objects
- Jets (*j*).
 - ▶ Energy deposit grouping with *infra*-red safe algorithm:
 - ▶ $p_T > 25$ GeV and $|\eta| < 4.5$
 - ♦ Clustering with anti- $k_{\rm T}$, R=0.4



G. Barone



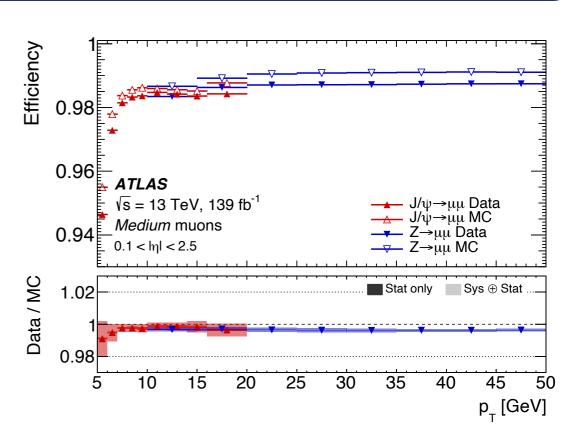


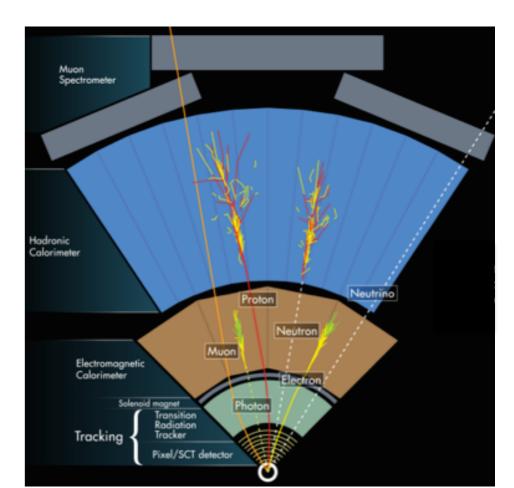
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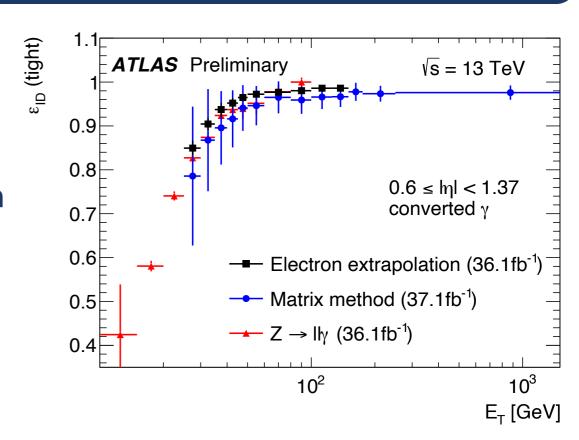
G. Barone

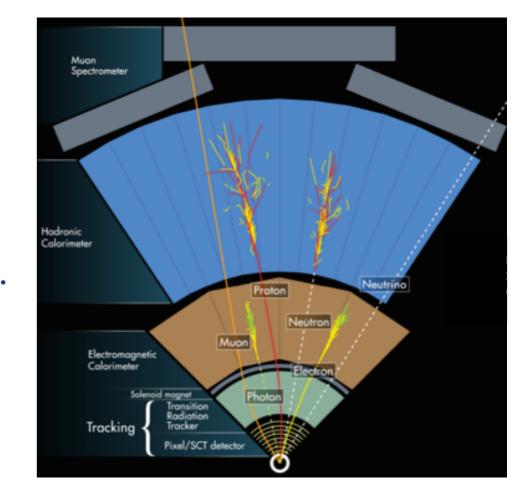




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 - Isolated objects
- Photons (γ) .
 - Clustering of calorimeter energy deposits.
 - ▶ Identified with rectangular cuts on shower shapes.



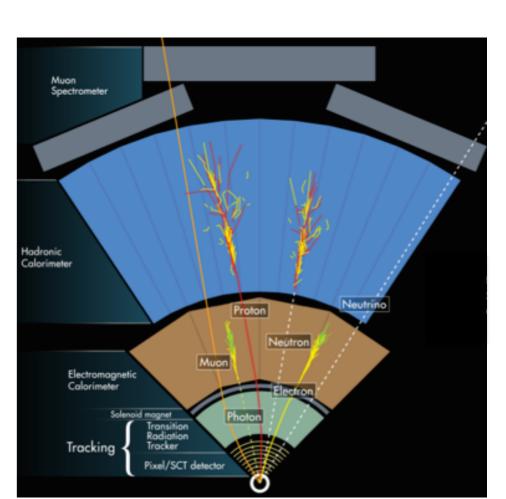




Reconstruction and selection

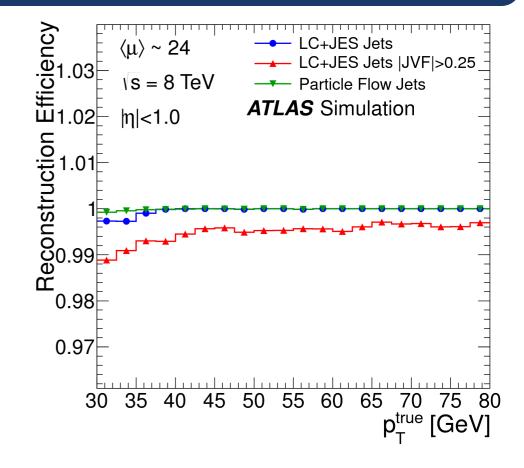
- Electrons (e).
 - Isolated objects clustered from calorimeter energy deposits with associated ID track.
 - ▶ $E_T > 7$ GeV, $|\eta| < 2.47$ and $|z_0 \sin(\vartheta)| < 0.5$ mm
- Muons (μ) .
 - Combined track fit of Inner Detector and Muon Spectrometer hits,
 - ▶ $p_T > 5$ GeV, $|\eta| < 2.7 |z_0 \sin(\vartheta)| < 0.5$ mm of "loose or medium quality"
 - Isolated objects
- Jets (*j*).
 - Combined from energy deposits in calorimeters and charged particles momenta
 - ▶ p_T > 20 GeV and $|\eta|$ < 4.5
 - ♦ Clustering with anti- $k_{\rm T}$, R=0.4

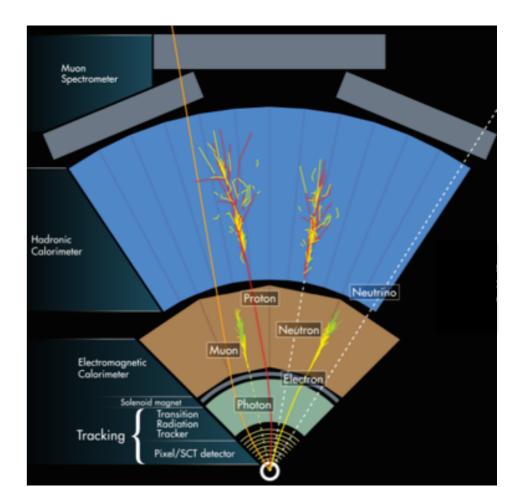




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 - Isolated objects
- Missing transverse energy ($E_{\rm T}^{\rm miss}$).
 - ▶ Inferred from transverse momentum imbalance







Couplings interpretations

- Interpretation of couplings cross sections in the context of new physics.
- Assuming production and decay are factorised

$$\sigma_i \times \mathbf{B}_f = \frac{\sigma_i(\kappa) \times \Gamma_f(\kappa)}{\Gamma_H}$$

 $\kappa_V = 1.03 \pm 0.03$

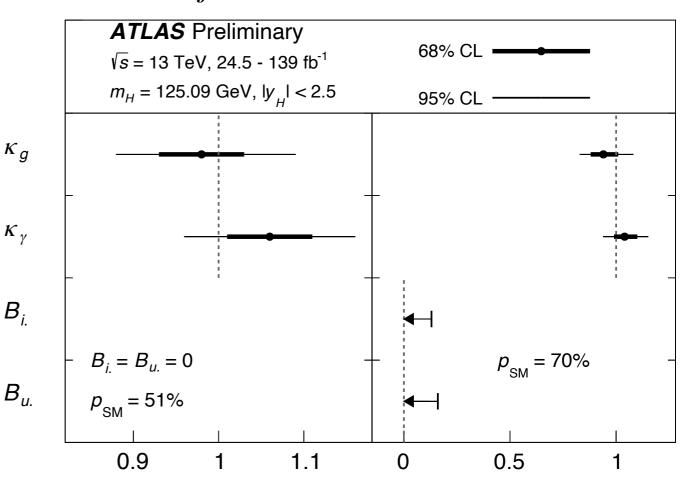
 $\kappa_F = 0.97 \pm 0.07$.

Coupling strength modifiers

$$\kappa_j^2 = \frac{\sigma_j}{\sigma_i^{\text{SM}}}$$

$$\kappa_j^2 = \frac{\Gamma_j}{\Gamma_i^{\text{SM}}}$$

- BSM contributions in loops and decays
 - k_g and k_Y measured with all other modifiers fixed to SM value.
 - Both hypotheses of invisible decays $B_{i.}$ (B_{inv} and B_{und} floating, with $k_F = k_V = 1$) and no invisible decays ($B_{inv} = B_{und} = 0$). $B_{u.}$



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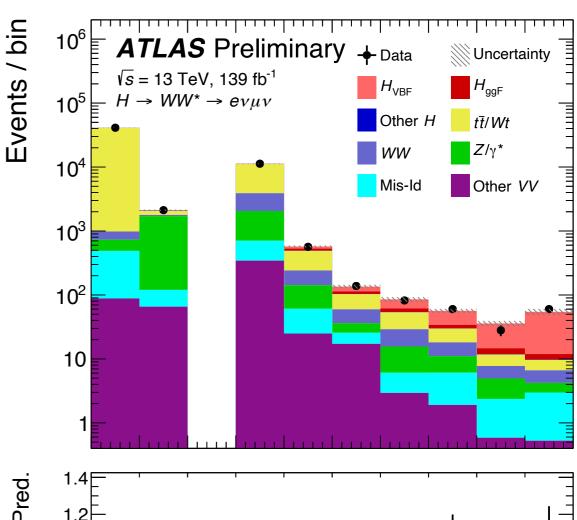
VBF in $WW^{\!*}$

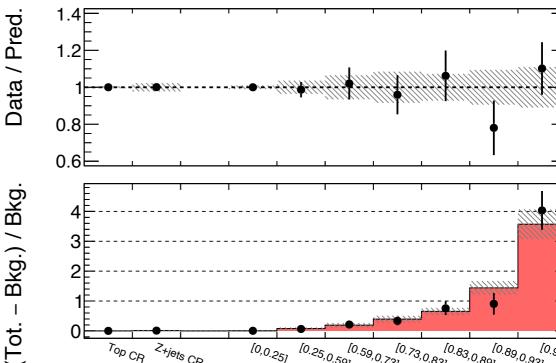
Production modes

- $WW^*{
 ightarrow}e
 u\mu
 u$ selection
 - ▶ Two isolated leptons $p_T(\ell) > 22$ GeV and $p_T(\ell) > 15 \text{ GeV}$
 - ▶ E_Tmiss > 20 GeV
 - Neural network for VBF production
- Neural network for VBF production
- Background estimation:

- I. Non resonant WW production
- 2. tt production
- 3. Drell-Yan: $Z \rightarrow \tau^+ \tau^-$
- 4. Hadrons misidentified as leptons:
 - W+jets $t\overline{t}$ and WZ production
- 5. ZZ^* , WZ, $W\gamma(*)$ production in
- 6. Single-top-quark (Wt) production

on simulation



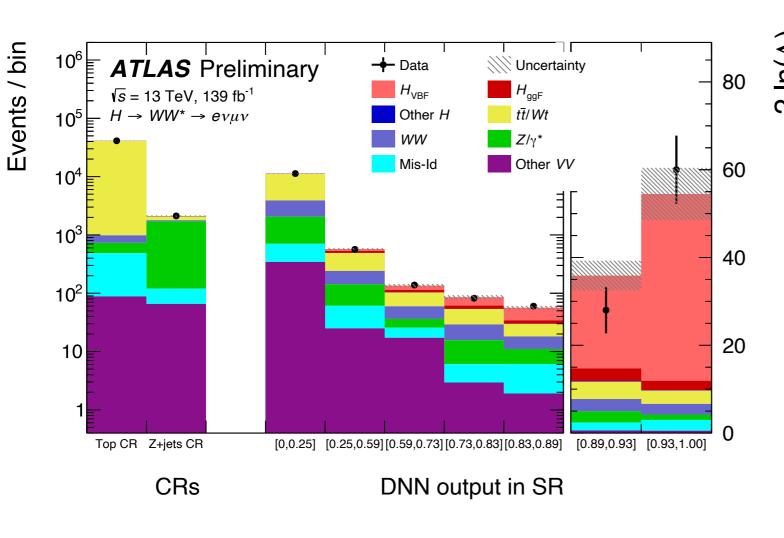


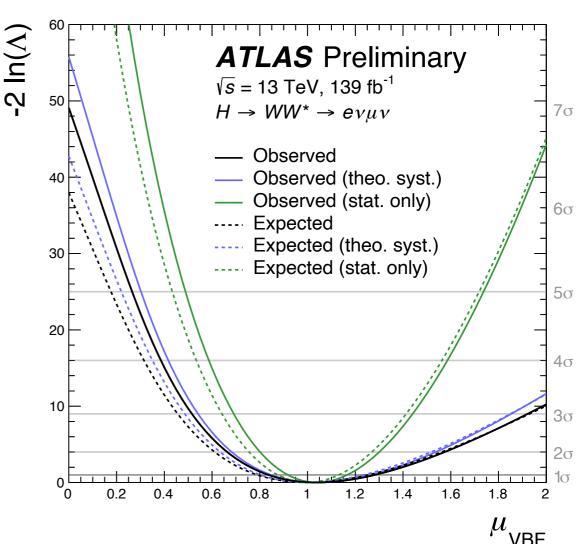
CRs

[0.59,0.73]

 $\frac{10.73,0.83j}{10.89,0.93j}$

DNN output in SR



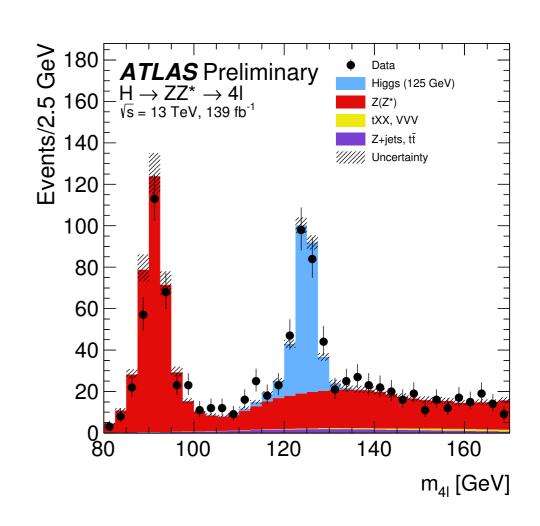


- ullet Measured production cross section $0.85\pm0.1(\mathrm{stat})$ $^{+0.17}$ - $_{0.13}$ (sys) pb
 - \blacktriangleright Significance of the signal observed of 7.0 σ
 - ♦ 6.2 σ expected for $m_{\rm H}$ = 125 GeV.



Reconstruction and selection

- $ZZ^* \rightarrow 4\ell \ (\ell = \mu,e)$ selection:
 - ▶ Isolated leptons with: $p_T(\ell) > 20 \text{ GeV}$, I5 GeV I0 GeV and 5 (7) GeV
 - \blacktriangleright Leading pair: pair closest to $m_{\mathbb{Z}}$,
 - Vertex refit: χ² cut at 99.5% signal efficiency
 - ▶ Final state photon emission recovered



- Background estimation
 - I. ZZ^* production in 4ℓ (dominant)
 - From $q\overline{q}$ annihilation and gg fusion (subdominant)
 - 2. Hadrons misidentified as leptons:
 - ightharpoonup Z+jets $t\overline{t}$ and WZ production
 - Extrapolation to signal region making use of simulation

3. ZZZ, WZZ and WWZ (small).

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Based on data

Based on simulation

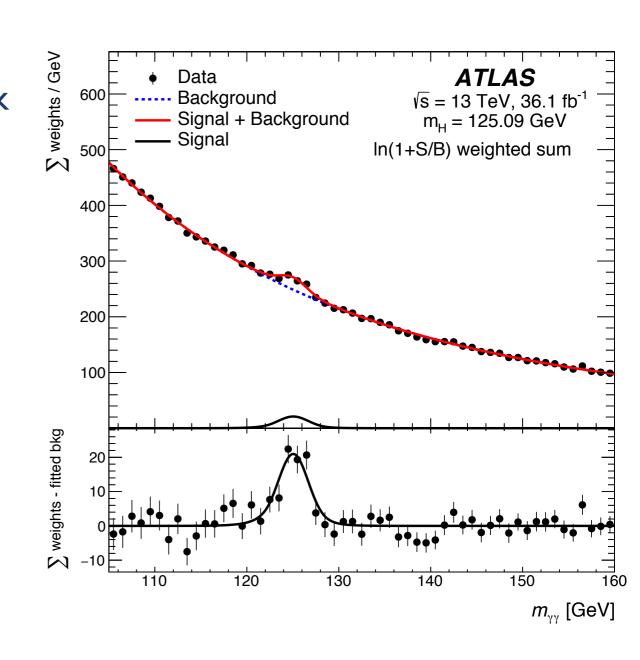
Reconstruction and selection

Diphoton event selection

- ▶ At least two photon with E_T > 25 GeV
- \blacktriangleright Highest E_T pair forms candidate.
- Vertex identification with Neural Network
 - ◆ Vertex within 0.3 mm for 79% of ggH events.

Background estimation

- ▶ Entirely estimated from data
- ▶ Prompt photons: maximum likelihood fit to m_{YY} spectrum
- Jets misidentified as photons: from control sample

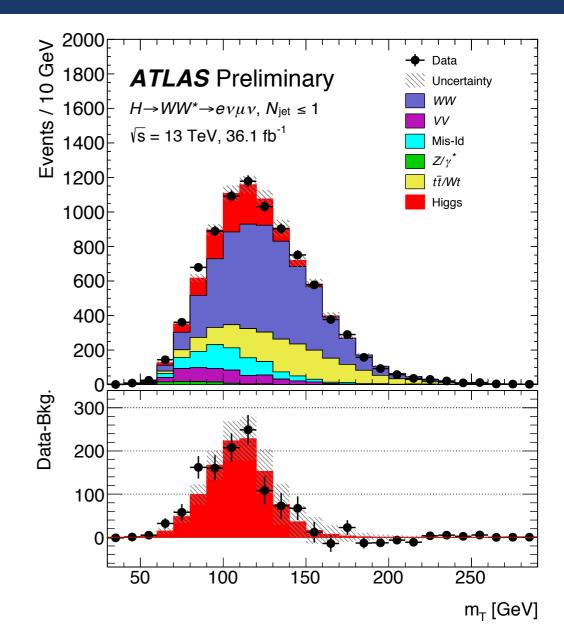


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$WW^* \rightarrow \ell \overline{\nu} \overline{\ell} \nu$

Reconstruction and selection

- $WW^* \rightarrow e \nu \mu \nu$ selection
 - ▶ Two isolated leptons $p_T(\ell) > 22$ GeV and $p_T(\ell) > 15$ GeV
 - ► E_Tmiss > 20 GeV
- Signal-to-background discriminants
 - ▶ Trasnverse mass (m_T) for ggF production and Boosted Decision Tree (BDT) for VBF production
- Background estimation
- I. Non resonant WW production
- 2. $t\overline{t}$ production
- 3. Drell-Yan: $Z{
 ightarrow} au^+ au^-$
- 4. Hadrons misidentified as leptons:
 - W+jets $t\overline{t}$ and WZ production



- 5. ZZ^* , WZ, $W\gamma(^*)$ production in
- 6. Single-top-quark (Wt) production

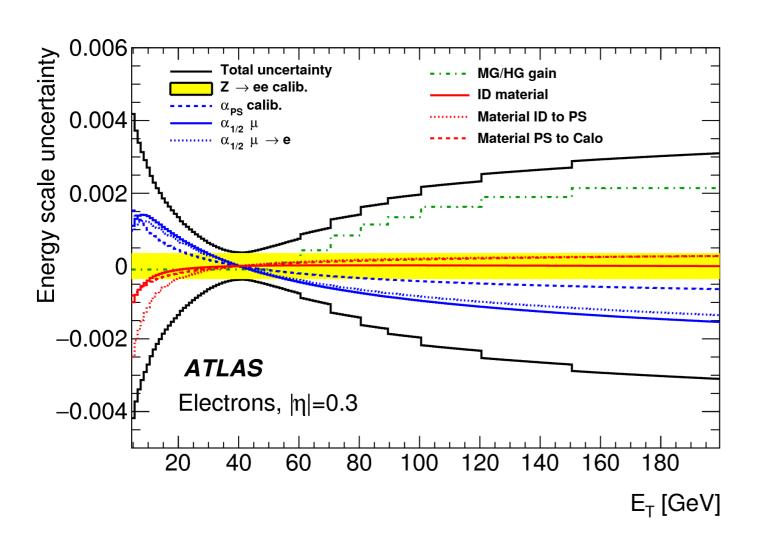
Based on data

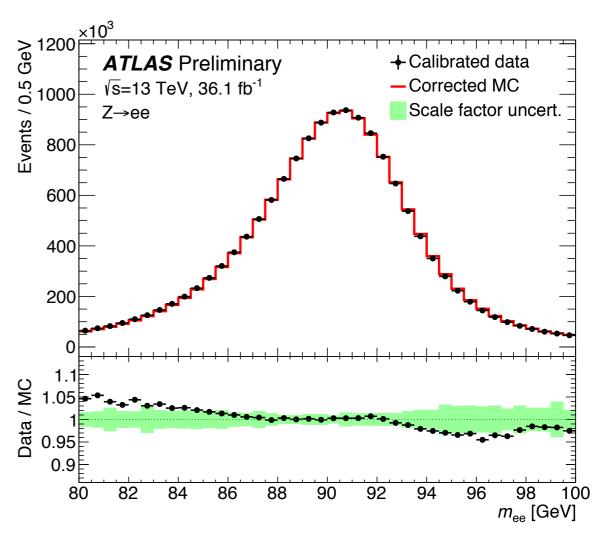
Based on simulation

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Mass measurement

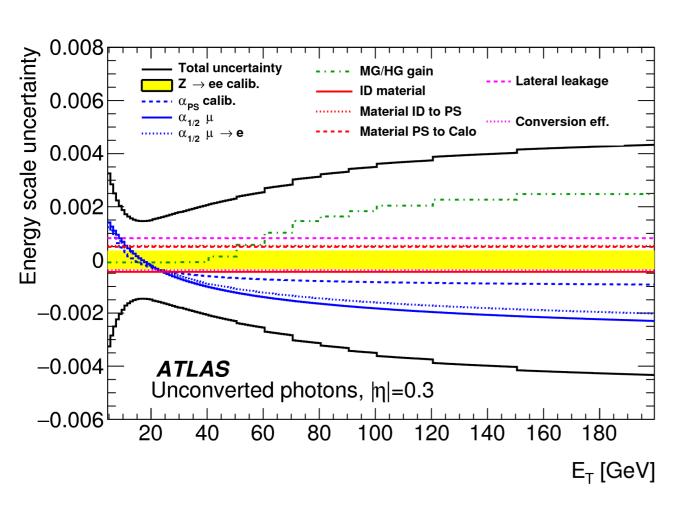
- \bullet Good energy calibration necessary for increased precision on m_H
 - ▶ Two step approach: i) material energy loss and ii) global calorimetric scale from $Z \rightarrow ee$ data
- Total scale uncertainty of at 40 GeV at the per-mille level.

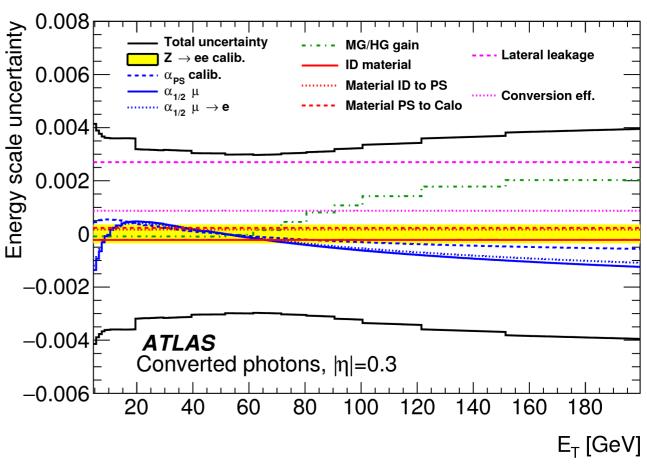




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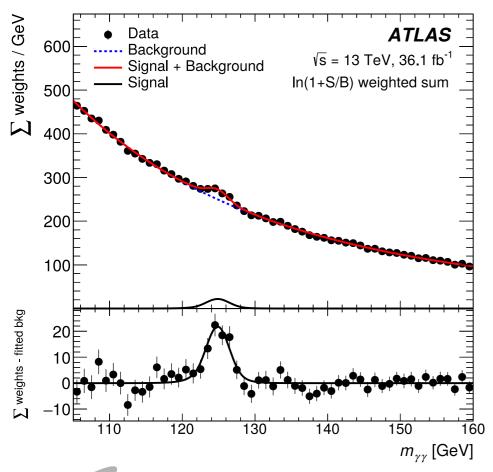
Mass measurement

• $H \rightarrow \gamma \gamma$ updated result at Run II.

arXiv:1806.00242

- Analytical function in kinematic and detector categories.
- Reduction of uncertainty through categorisation of events as a function of resolution and signal significance.
- Expected statistical uncertainty of 0.21 GeV and 0.34 GeV systematic uncertainty

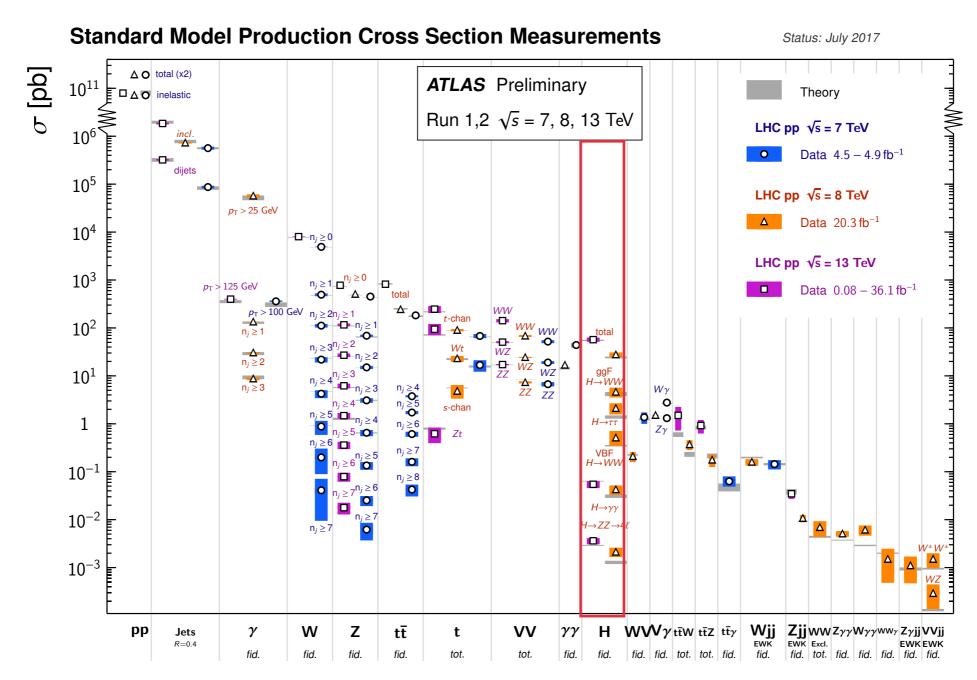
$$m_H^{\gamma\gamma} = 124.93 \pm 0.40 \ (\pm 0.21 \ \mathrm{stat} \ \mathrm{only}) \ \mathrm{GeV}$$



Source	Systematic uncertainty in m_H [MeV]
EM calorimeter response linearity	60
Non-ID material	55
EM calorimeter layer intercalibration	55
$Z \to ee$ calibration	45
ID material	45
Lateral shower shape	40
Muon momentum scale	20
Conversion reconstruction	20
$H \to \gamma \gamma$ background modelling	20
$H \to \gamma \gamma$ vertex reconstruction	15
e/γ energy resolution	15
All other systematic uncertainties	10

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Introduction



- Higgs boson decay channels considered in this talk :
 - (i) $H \rightarrow Dibosons (ZZ^* \rightarrow 4\boldsymbol{\ell}, WW^* \rightarrow \boldsymbol{\ell} \overline{\nu} \overline{\boldsymbol{\ell}} \nu)$, and $\gamma \gamma$
 - (ii) $H \rightarrow \text{fermions } (\tau \overline{\tau}, b \overline{b}, \mu \overline{\mu})$



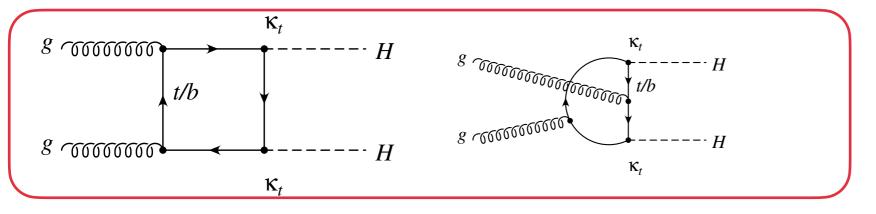
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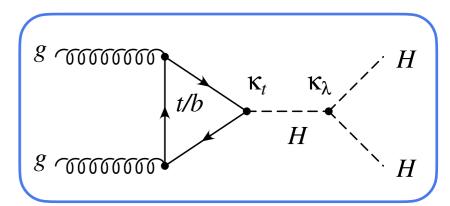
Di-Higgs production

 The Higgs boson can interact with itself through quadratic terms in the Higgs potential

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i \bar{\Psi} D^{\mu} \Phi + V(\Phi) \qquad V(\Phi) \sim -\mu^{2} (\Phi^{\dagger} \Phi) + \lambda (\Phi^{\dagger} \Phi)^{2} + \bar{\Psi}_{L} \hat{Y} \Phi \Psi_{R} + h.c.$$

- ▶ About 500 times suppression of $\sigma(gg \rightarrow H)$ (48.5 pb) $/ \sigma(gg \rightarrow HH)$ (~33.4 fb)
- Destructive interference between the terms proportional to the K_t^2 and the product of K_t and K_{λ}





- Single Higgs process do not depend on the trilinear self (λ_{HHH}) coupling at LO but are needed for the NLO EW corrections.
- Indirect constraint with comparing NLO EW dependant λ_{HHH} effects

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Di-Higgs production

Modified Higgs production cross section and branching rations to account

for NLO EW corrections (K^{i}_{EW} and C^{f})

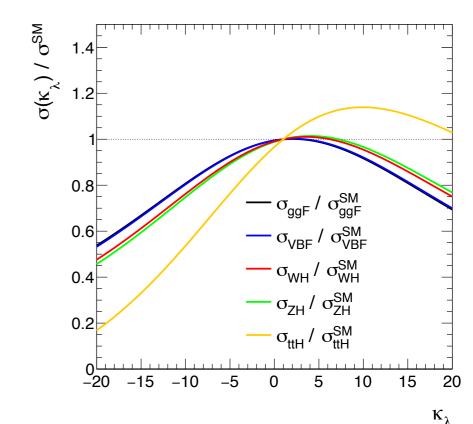
$$\mu_{i}(\kappa_{\lambda}, \kappa_{i}) = \frac{\sigma^{\text{BSM}}}{\sigma^{\text{SM}}} = Z_{H}^{\text{BSM}}(\kappa_{\lambda}) \left[\kappa_{i}^{2} + \frac{(\kappa_{\lambda} - 1)C_{1}^{i}}{K_{\text{EW}}^{i}} \right]$$

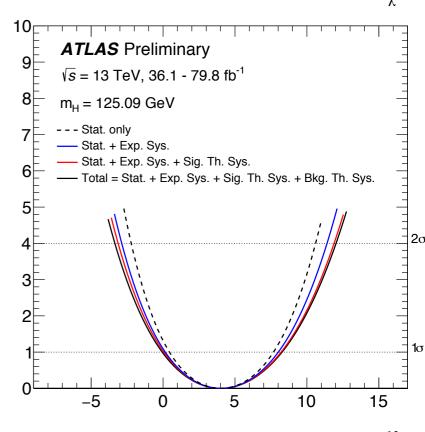
$$\mu_f(\kappa_{\lambda}, \kappa_f) = \frac{\mathrm{BR}_f^{\mathrm{BSM}}}{\mathrm{BR}_f^{SM}} = \frac{\kappa_f^2 + (\kappa_{\lambda} - 1)C_1^f}{\sum_j \mathrm{BR}_j^{\mathrm{SM}} \left[\kappa_j^2 + (\kappa_{\lambda} - 1)C_1^j\right]}$$

Combined fit over Higgs STXS combination

$$\kappa_{\lambda} = 4.0^{+4.3}_{-4.1} = 4.0^{+3.7}_{-3.6} (\text{stat.})^{+1.6}_{-1.5} (\text{exp.})^{+1.3}_{-0.9} (\text{sig. th.})^{+0.8}_{-0.9} (\text{bkg. th.})$$

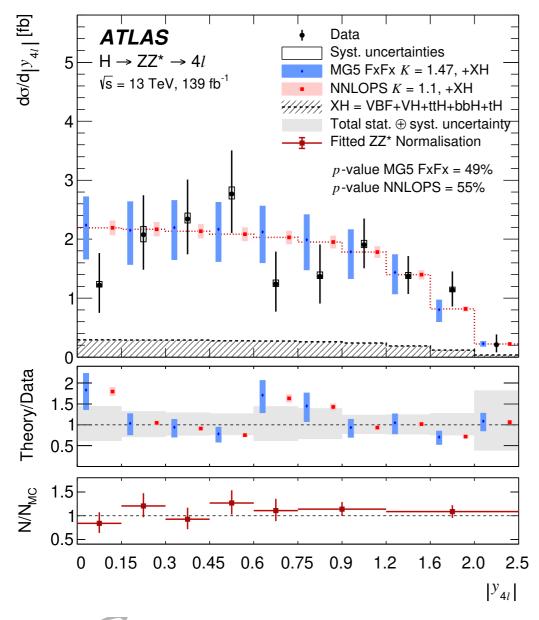
▶ 95% C.L. -3.2 < κ_{λ} < 11.9 comparable over direct HH searches (-5.0 < κ_{λ} < 12.1)

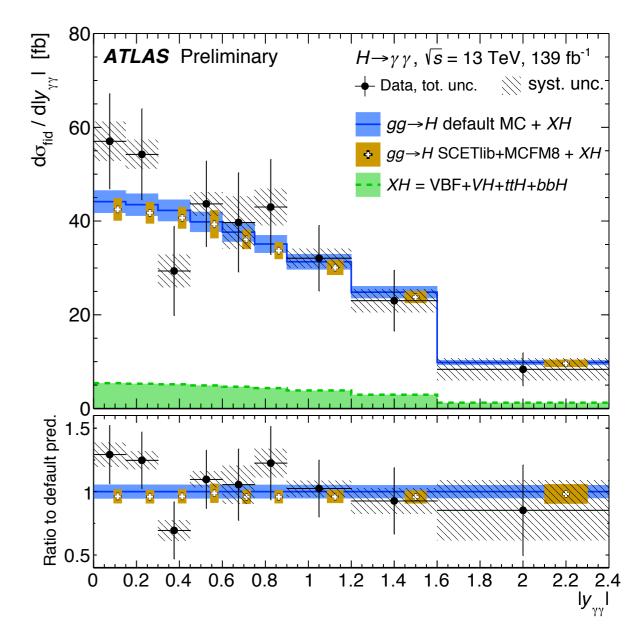




Higgs boson kinematics

- 4ℓ : isolate signal under the Higgs resonant peak (115 < $m_{4\ell}$ < 130).
- $\gamma\gamma$ cross section extracted from resonant peak over the $\gamma\gamma$ continuum.
- Higgs boson $p_{T,4\ell(\gamma\gamma)}$ and rapidity $(y_{4\ell(\gamma\gamma)})$ probe.
 - $p_{T,4\ell(\gamma\gamma)}$: Lagrangian structure of H interactions, Yukawa couplings
 - $y_{4\ell(\gamma\gamma)}$: Sensitivity to proton's parton density functions.





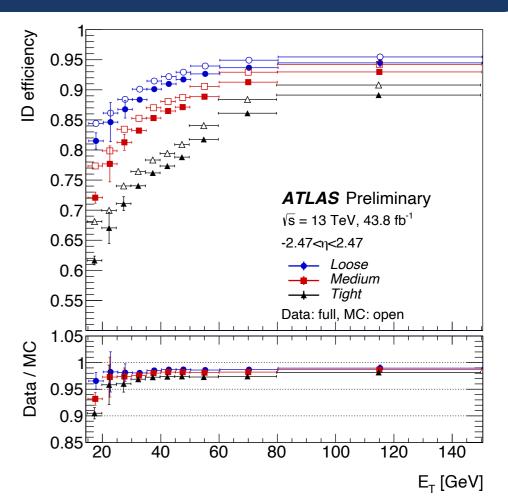
Physics objects definition and selection criteria.

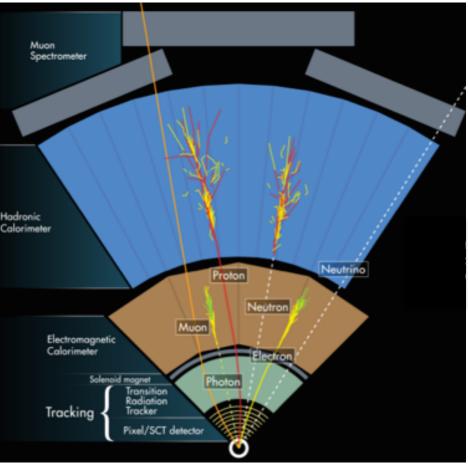
Reconstruction and selection

- Electrons (e).
 - Isolated objects clustered from calorimeter energy deposits with associated ID track.
 - ▶ $E_T > 7$ GeV, $|\eta| < 2.47$ and $|z_0 \sin(\vartheta)| < 0.5$ mm
- Muons (μ) .
 - Combined track fit of Inner Detector and Muon Spectrometer hits,
 - ▶ $p_T > 5$ GeV, $|\eta| < 2.7 |z_0 \sin(\vartheta)| < 0.5$ mm of "loose or medium quality"
 - Isolated objects
- Jets (*j*).
 - ▶ Energy deposit grouping with *infra*-red safe algorithm:
 - ▶ $p_T > 20$ GeV and $|\eta| < 4.5$
 - ♦ Clustering with anti- $k_{\rm T}$, R=0.4



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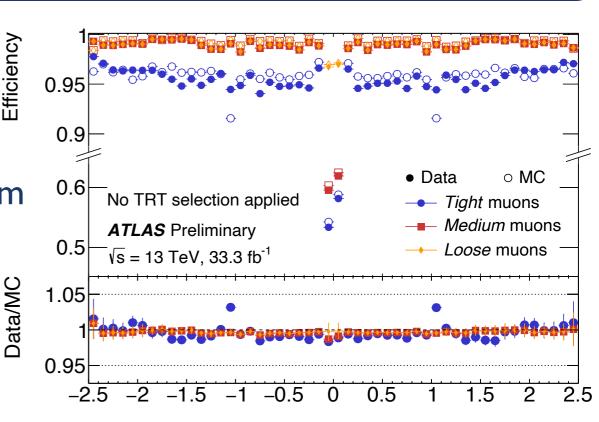


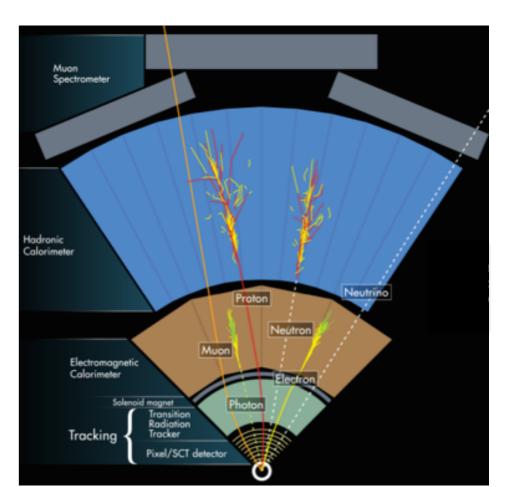
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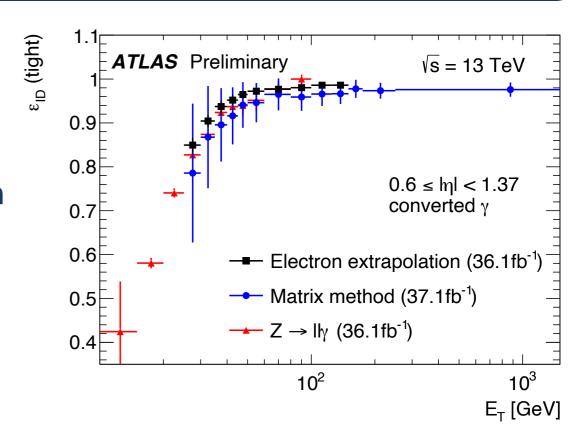
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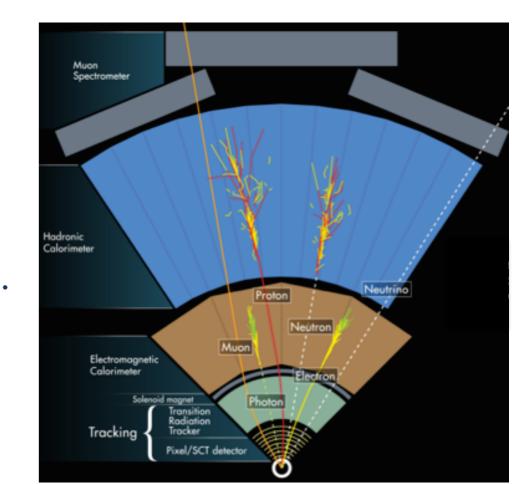




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- Photons (γ) .
 - Clustering of calorimeter energy deposits.
 - ▶ Identified with rectangular cuts on shower shapes.





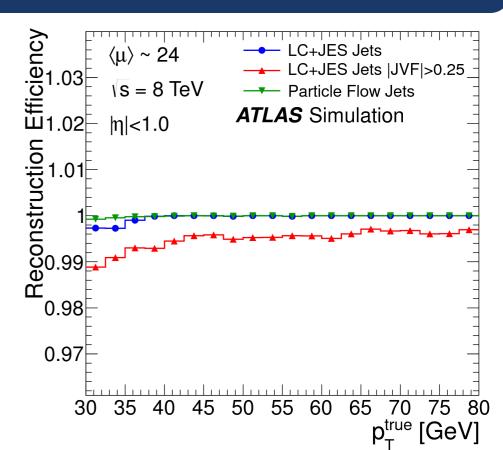


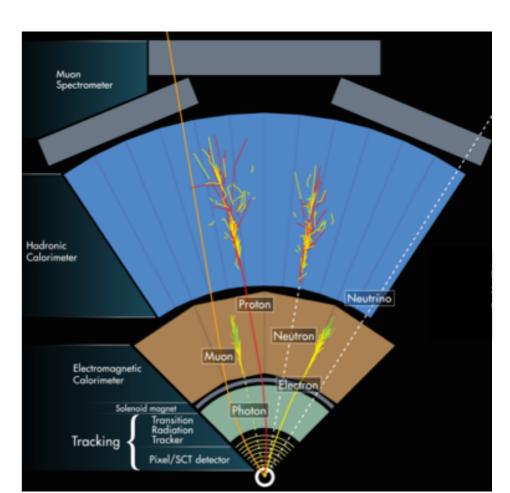
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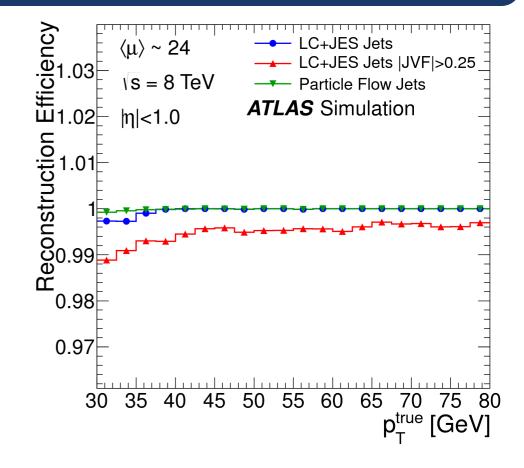
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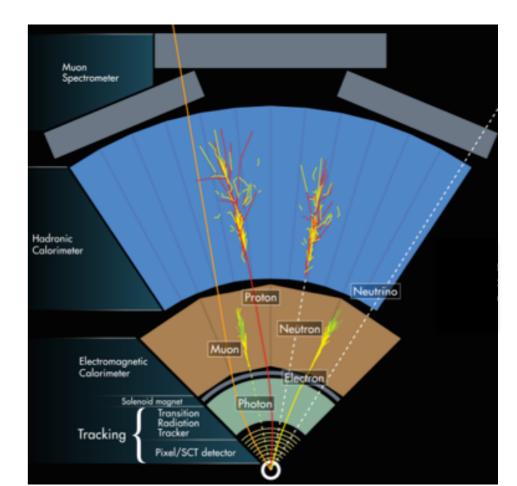




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 - ▶ Inferred from transverse momentum imbalance







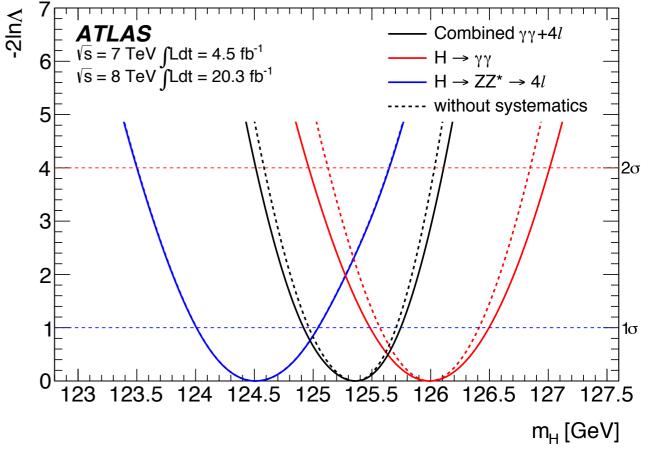
Run I status

- ATLAS run I precision on m_H of 0.33%
 - combined measurement from $H \rightarrow \gamma \gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$.

Channel	Mass measurement [GeV]		
$H o \gamma \gamma$	$125.98 \pm 0.42 \text{ (stat)} \pm 0.28 \text{ (syst)} = 125.98 \pm 0.50$		
$H \to ZZ^* \to 4\ell$	$124.51 \pm 0.52 \text{ (stat)} \pm 0.06 \text{ (syst)} = 124.51 \pm 0.52$		
Combined	$125.36 \pm 0.37 \text{ (stat)} \pm 0.18 \text{ (syst)} = 125.36 \pm 0.41$		

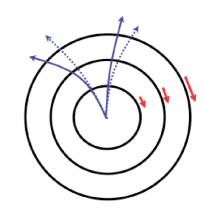
▶ For both channels dominated by statistical uncertainty

- Aim in improving significantly on δm_H
 - ▶ Expect 1.7 times more candidates, with 36 fb⁻¹ at \sqrt{s} =13 TeV



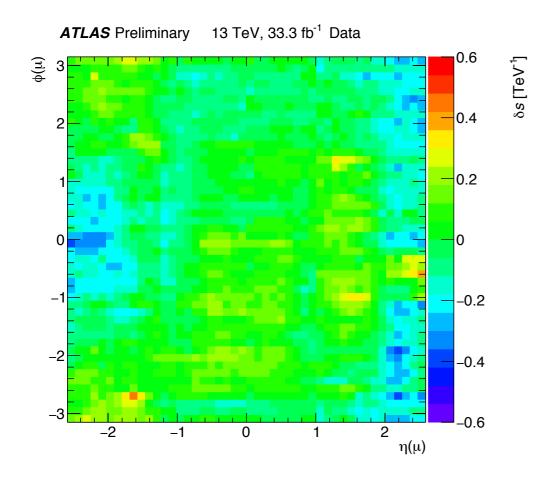
Mass Measurement

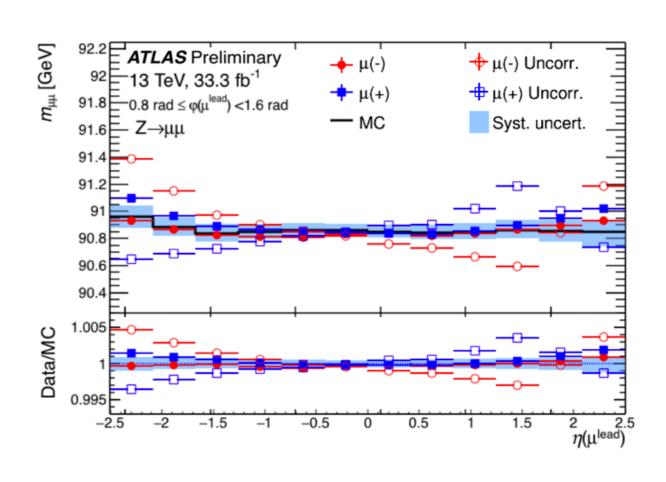
- Correction for local misalignments
 - ▶ Charge dependent bias, with net effect of worsening resolution
 - In-situ correction based on $Z \rightarrow \mu\mu$ data, recovers up to 5% in resolution.



• Iteratively removing the bias δ_s :

$$p_{\mathrm{T}}^{\mathrm{corr}}(\mu) = \frac{p_{\mathrm{T}}^{\mathrm{bias}}(\mu)}{1 - q(\mu)\delta_{s}(\eta, \phi)p_{\mathrm{T}}^{\mathrm{bias}}(\mu)}$$



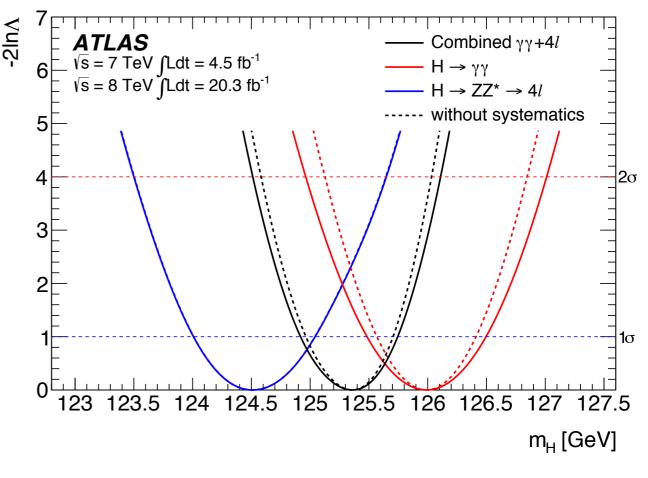


Mass measurement

- ATLAS run I precision on m_H of 0.33%
 - combined measurement from $H \rightarrow \gamma \gamma$ and $H \rightarrow ZZ^* \rightarrow 4\boldsymbol{\ell}$.

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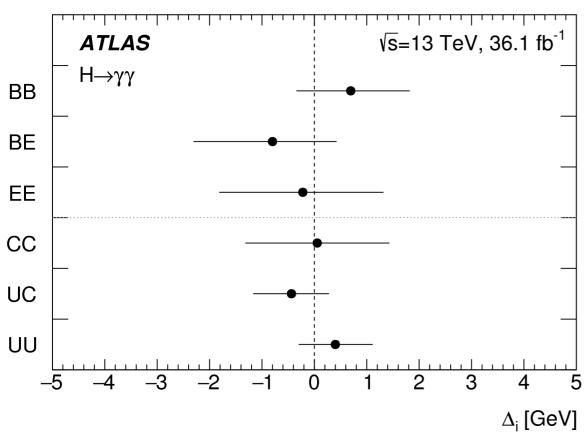
- ▶ For both channels dominated by statistical uncertainty
- ullet Compatibility within 2.0 σ
 - ▶ p-value of about 0.05.
- Aim in improving significantly on δm_H
 - ▶ Expect 1.7 times more candidates, with 36 fb⁻¹ at \sqrt{s} =13 TeV

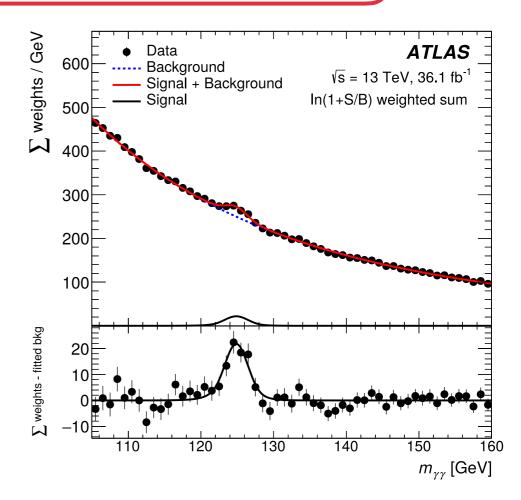


arXiv:1806.00242

- $H \rightarrow \gamma \gamma$ updated result at Run II.
 - Analytical function in kinematic and detector categories.
 - Reduction of uncertainty through categorisation of events as a function of resolution and signal significance.
- Expected statistical uncertainty of 0.21 GeV and 0.34 GeV systematic uncertainty

$$m_H^{\gamma\gamma} = 124.93 \pm 0.40 \ (\pm 0.21 \ {\rm stat \ only}) \ {\rm GeV}$$







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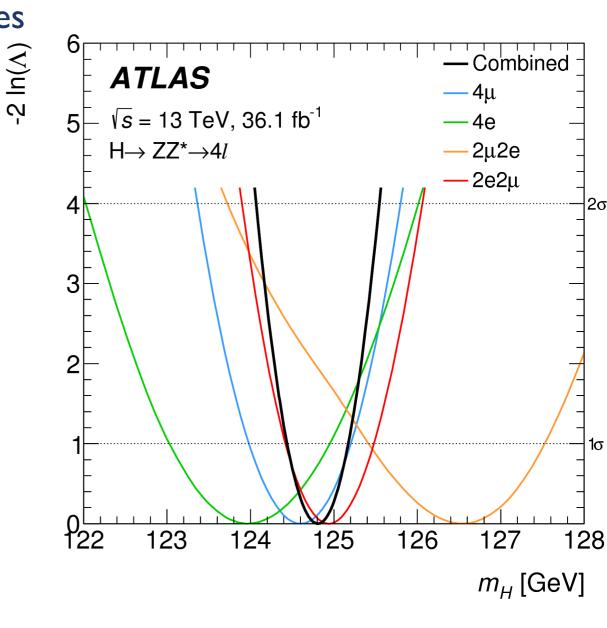
January-21

arXiv:1806.00242

- Final estimate from 4x4 simultaneous un-binned fit
 - ▶ Four kinematic categories and four final states
- Good agreement between channels.
- Systematic uncertainty of 50 MeV

Systematic effect	Uncertainty on $m_H^{ZZ^*}$ [MeV]
Muon momentum scale	40
Electron energy scale	26
Pile-up simulation	10
Simulation statistics	8

• Result:



▶ 25% improved precision with respect to Run I ATLAS Combination.

$$m_H^{ZZ^*} = 124.79 \pm 0.36 \ (\pm 0.05 \ \mathrm{stat} \ \mathrm{only}) \ \mathrm{GeV}$$



Reconstruction and selection

ATLAS-CONF-2018-018

- $ZZ^* \rightarrow 4\ell$ ($\ell = \mu,e$) selection:
 - ▶ Isolated leptons with: $p_T(\ell) > 20 \text{ GeV}$, I5 GeV I0 GeV and 5 (7) GeV
 - \blacktriangleright Leading pair: pair closest to m_{Z_i}
 - Vertex refit: χ² cut at 99.5% signal efficiency
 - ▶ Final state photon emission recovered

Final	Signal	ZZ^*	Other	Total	Observed
state		background	backgrounds	expected	
$\overline{4\mu}$	40.5 ± 1.7	19.0 ± 1.1	1.71 ± 0.10	61.2 ± 2.0	64
$2e2\mu$	28.2 ± 1.2	13.3 ± 0.8	1.38 ± 0.10	42.8 ± 1.4	64
$2\mu 2e$	22.1 ± 1.4	9.2 ± 0.9	2.99 ± 0.09	34.3 ± 1.7	39
4e	21.1 ± 1.4	8.6 ± 0.8	2.90 ± 0.09	32.5 ± 1.6	28
Total	112 ± 5	50 ± 4	8.96 ± 0.12	171 ± 6	195

$$115~{
m GeV} < m_{4\it e} < 130~{
m GeV}$$

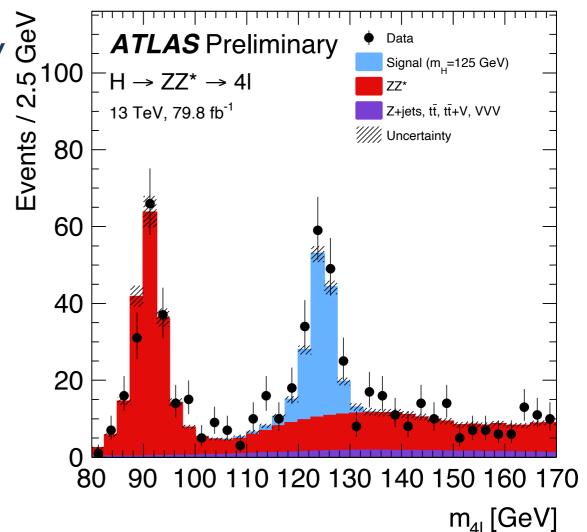
Background estimation

Based on simulation

- I. ZZ^* production in 4ℓ (dominant)
 - From $q\overline{q}$ annihilation and gg fusion (subdominant)

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2. ZZZ, WZZ and WWZ (small).



3. Hadrons misidentified as leptons:

- \blacktriangleright Z+jets $t\overline{t}$ and WZ production
- Extrapolation to signal region making use of simulation

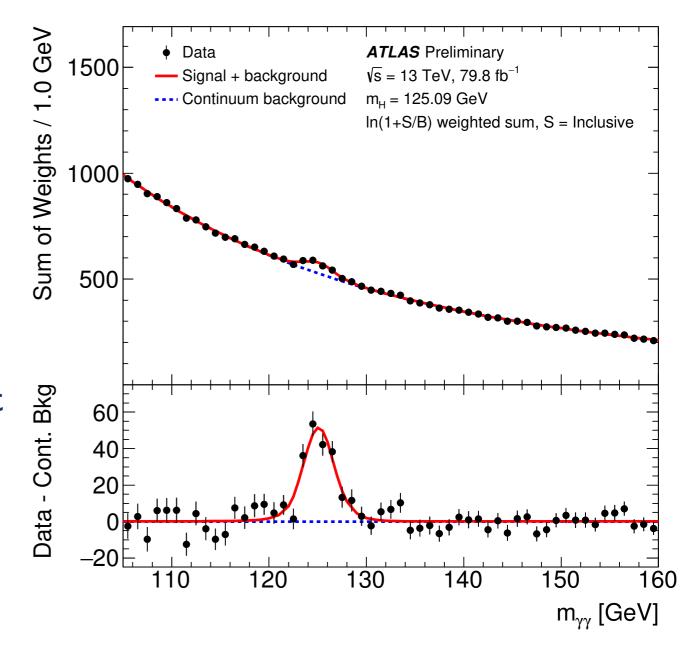
Based on data

ATLAS-CONF-2018-028

Diphoton event selection

- ▶ At least two photon with E_T > 25 GeV
- ▶ Highest E_T pair forms candidate.
- Vertex identification with Neural Network
 - ◆ Vertex within 0.3 mm for 79% of ggH events.

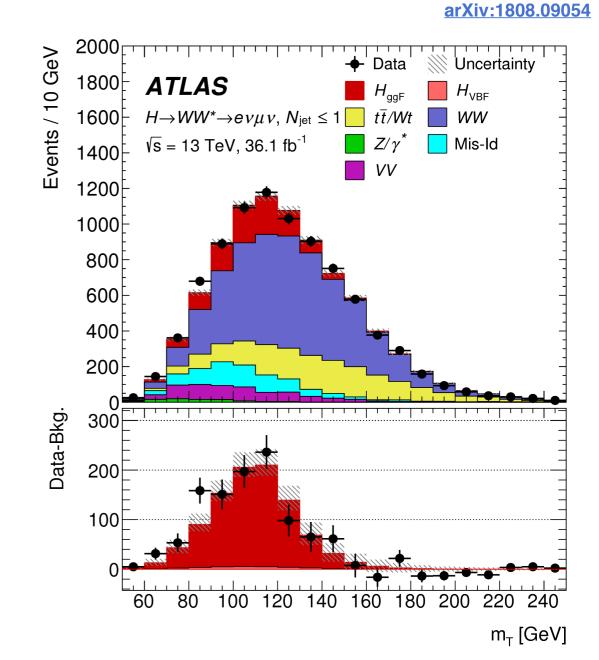
- Background estimation
 - ▶ Entirely estimated from data
 - ▶ Prompt photons: maximum likelihood fit to m_{YY} spectrum
 - Jets misidentified as photons: from control sample



$WW^* \rightarrow \ell \overline{\nu} \overline{\ell} \nu$

Reconstruction and selection

- $WW^* \rightarrow e \nu \mu \nu$ selection
 - ▶ Two isolated leptons $p_T(\ell) > 22$ GeV and $p_T(\ell) > 15$ GeV
 - ▶ E_Tmiss > 20 GeV
- Signal-to-background discriminants
 - ▶ Trasnverse mass (m_T) for ggF production and neural network for VBF production
- Background estimation Based on data
- 1. Non resonant WW production
- 2. $t\overline{t}$ production
- 3. Drell-Yan: $Z \rightarrow \tau^+ \tau^-$
- 4. Hadrons misidentified as leptons:
 - W+jets $t\overline{t}$ and WZ production



- 5. ZZ^* , WZ, $W\gamma(^*)$ production in
- 6. Single-top-quark (Wt) production

Based on simulation

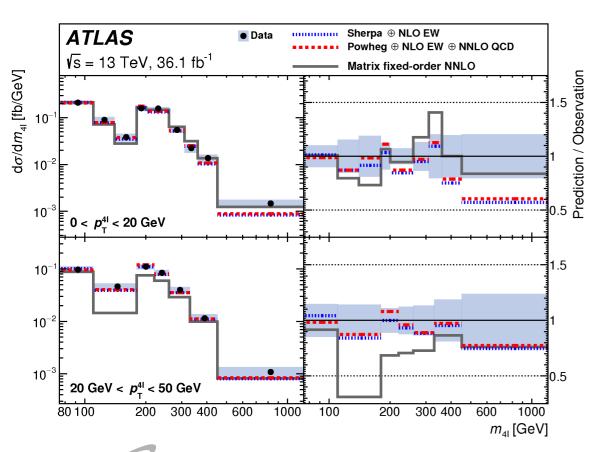
Production mode

- Cut based classification of events into category.
 - ▶ Ex. Jet multiplicity (ggF), m_{jj} for (VBF) and b-tagging $(t\overline{t}H)$
- and multivariate analysis (BDT) to discriminate contributions.
 - ▶ ggF from ZZ^* , VBF from ggF, VH(had) from all.
 - ▶ Variables: $p_{T,4\varrho}$, KD, η_j , $\Delta \eta_{jj}$, $p_{T,j}$ etc.
- Detector and theoretical uncertainties
- (i) Luminosity 3.2%
- (ii) Lepton Identification < 2%
- (iii)Pileup ~2%
- (iv)Jet Energy Scale (3%-7%)
- (v) Jet Energy Resolution (2%-4%)

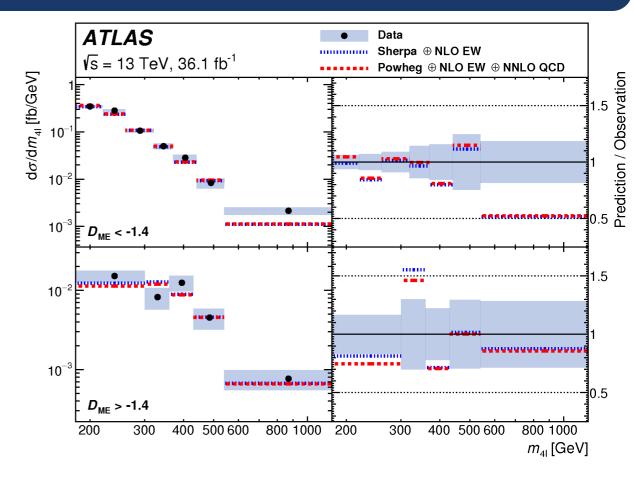
- (i) μ_R and μ_F about 4% to 30%
- (ii) ggF prediction in N_J categories.
- (iii)(BSM only NLO/LO prediction)

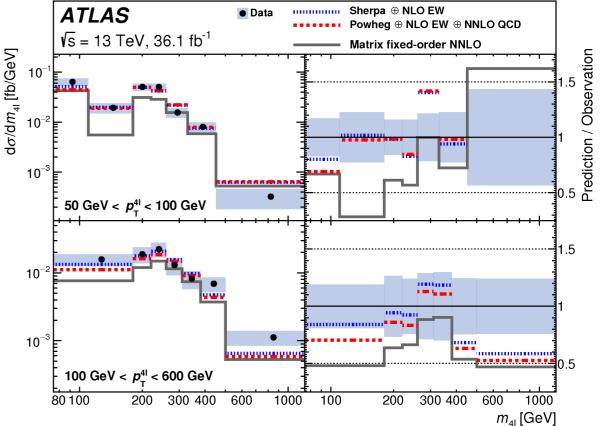
4*l* mass spectrum

- Double differential measurements as a function of rapidity and p_T
- ullet As well as matrix element discriminant between ZZ and H
 - ▶ MCFM-based.
- Comparisons with NLO EW and NLO EW and NNLO QCD predictions.



Differential cross section

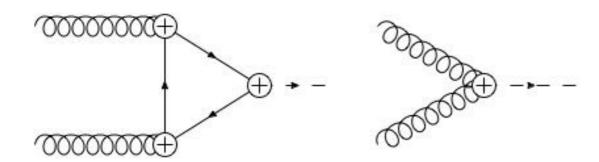




Higgs boson kinematics

Introduction

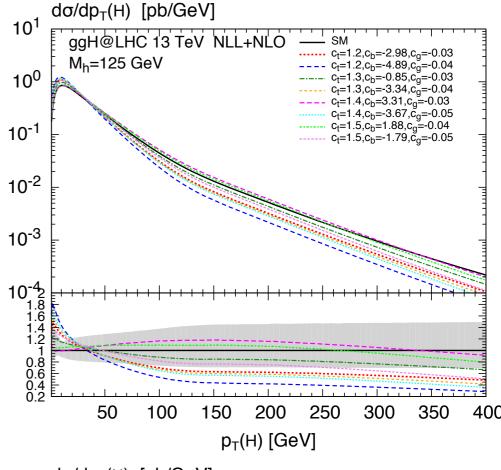
- Higgs boson $p_{T,4\ell}$ and rapidity $(y_{4\ell})$ probe:
 - ▶ $p_{T,4\ell}$: Lagrangian structure of H interactions.

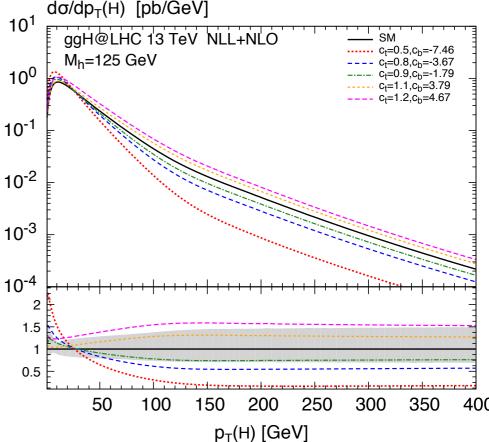


 Small perturbations to SM: dimension 6 operators most effective approach.

$$\begin{array}{c} \frac{c_1}{\Lambda^2}\,\mathcal{O}_1 \to \frac{\alpha_{\rm S}}{\pi v}c_g h G^a_{\mu\nu}G^{a,\mu\nu}\,, \\ \frac{c_2}{\Lambda^2}\,\mathcal{O}_2 \to \frac{m_t}{v}c_t h \bar{t}t\,, \\ \frac{c_3}{\Lambda^2}\,\mathcal{O}_3 \to \frac{m_b}{v}c_b h \bar{b}b\,, \\ \end{array} \qquad \begin{array}{c} \mathbf{c}_{\rm t}: t \text{ and } b \text{ Yukawa couplings} \\ \frac{c_4}{\Lambda^2}\,\mathcal{O}_4 \to c_{tg} \frac{g_S m_t}{2 v^3} (v+h) G^a_{\mu\nu} (\bar{t}_L \sigma^{\mu\nu} T^a t_R + h.c) \end{array}$$

 c_{tg} : dipole-moment, g-t interaction





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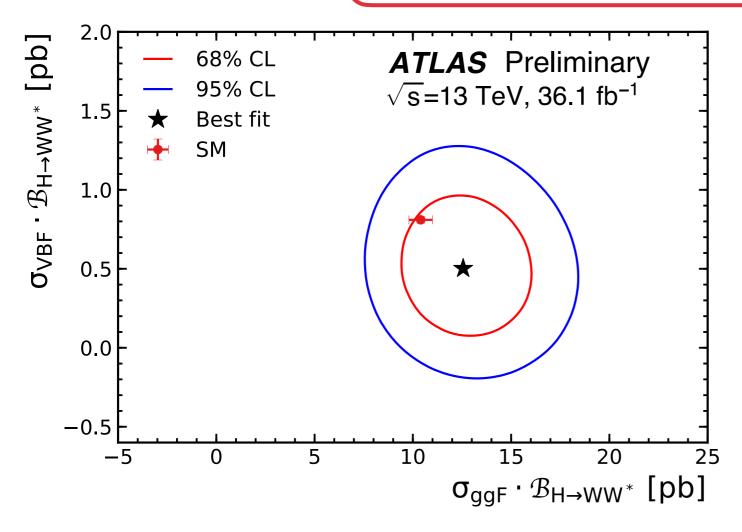
Grazzini et al. arXiv:1612.00283 January-21

$WW^* \rightarrow \ell \overline{\nu} \overline{\ell} \nu$ results

- Simultaneous fit to the ggF and VBF categories.
 - ▶ Over m_T for ggF and BDT response for ggF
 - ▶ Extraction of ggF and VBF total cross sections

$$\mu_{\text{ggF}} = 1.21^{+0.12}_{-0.11}(\text{stat.})^{+0.18}_{-0.17}(\text{sys.})$$

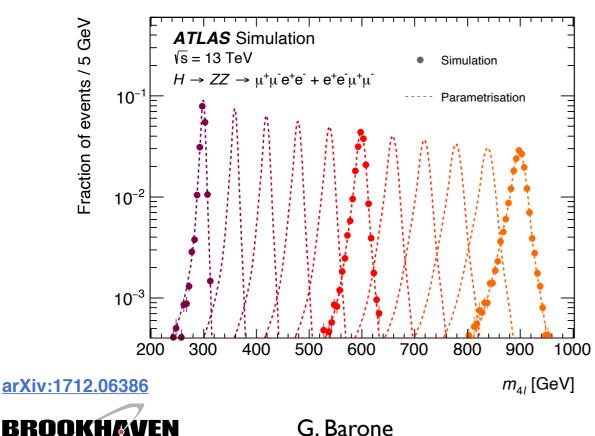
$$\mu_{\text{VBF}} = 0.62^{+0.30}_{-0.28}(\text{stat.}) \pm 0.22(\text{sys.})$$

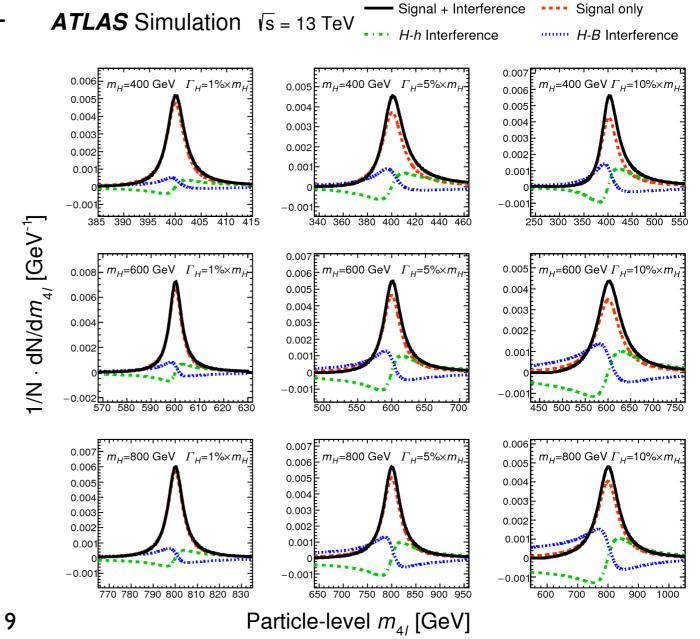


ullet Significances of 6.3 σ and 1.8 σ for ggF and VBF, respectively

$ZZ \rightarrow 4\ell$ and $\ell \overline{\ell} \nu \overline{\nu}$

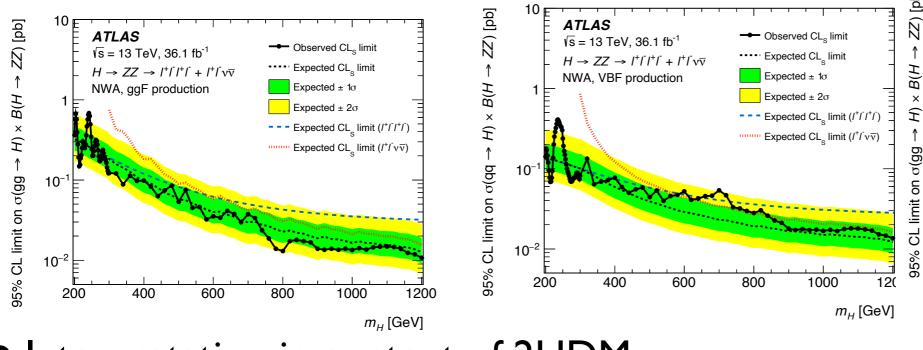
- Searches for spin-0 and spin-2 resonances in the $ZZ \rightarrow 4\ell$ and $\ell \overline{\ell} \nu \overline{\nu}$ final states.
 - Upper limits for Type-I and II two-Higgs double models (spin-0) and for RS models (spin-2)
 - Separate sensitivity for ggF and VBF productions (both ATLAS and CMS)
 - ♦ Typical VBF selection: at least two jets with $p_T(j) > 30$ GeV, $\Delta \eta > 3.3$ and $m_{jj} > 400$ GeV
- Resonances searched in $m_{4\ell}$ and m_T
 - Analytical parametrisation of signal.
 - ▶ h-H interference taken into account in the large width approximation

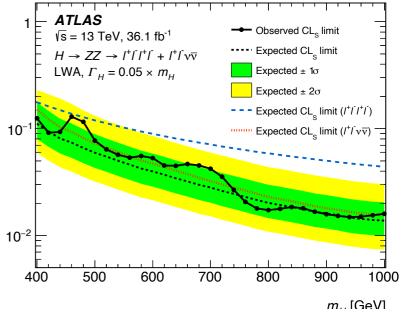




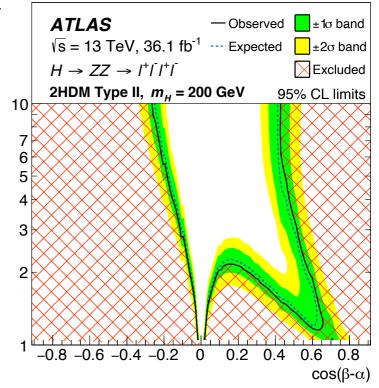
$ZZ \rightarrow 4\ell$ and $\ell \overline{\ell} \nu \overline{\nu}$

- Searches for spin-0 and spin-2 resonances in the $ZZ(\to 4\ell)$ and $\ell \overline{\ell} \nu \overline{\nu}$ channel
- Spin-0 resonance limits
 - Narrow width: 0.68 pb at m_H = 242 GeV to 11 fb at m_H = 1.2 TeV
 - ▶ Large width as a function of 1%, 5% and 10% of m_H





- Interpretation in context of 2HDM
 - No direct coupling of Higgs to leptons, only Type II and I considered.
 - Relative ggF to VBF rates fixed to 2HDM predictions for m_H = 200 GeV.
 - ♦ NWA valid across wide range and maximal experimental sensitivity
 arXiv:1712.06386



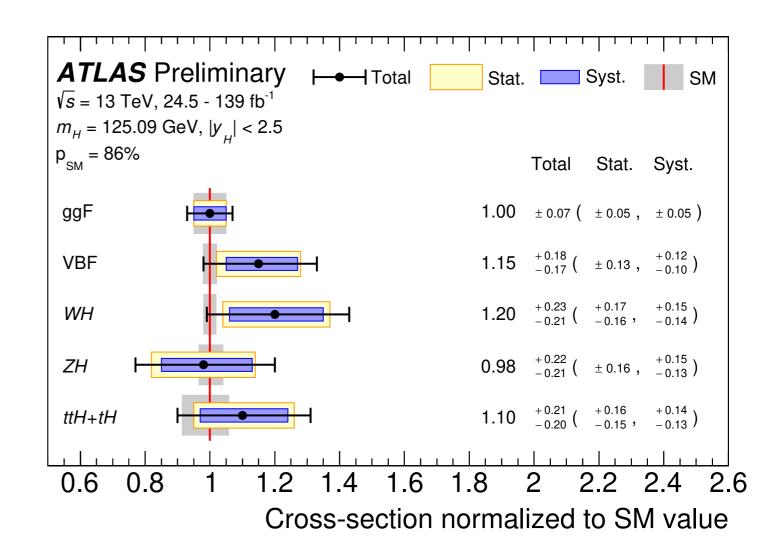
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Production mode

- Improved statistical sensitivity
 - reduced correlations between production modes.
- Statistical and systematic uncertainties of the same size.



Process Value Uncertainty [pb] SM pred. $(|y_H| < 2.5)$ Total + Exp. Sig. Th. [pb]Stat. Syst. Bkg. Th. [pb]+ 0.9+ 1.0ggF 44.7 ± 3.1 ± 2.2 44.7 ± 2.2 -1.7-0.9-0.7+ 0.3 $3.51 \, {}^{+\ 0.08}_{-\ 0.07}$ **VBF** ± 0.6 ± 0.1 4.0 ± 0.5 ± 0.4 ± 0.3 $^{+\ 0.28}_{-\ 0.25}$ $^{+\ 0.20}_{-\ 0.19}$ + 0.08 $+\ 0.10 \\ -\ 0.09$ + 0.13WH1.45 1.204 ± 0.024 -0.06-0.12 $+ 0.18 \\ - 0.17$ $0.797 \, {}^{+}_{-}\, {}^{0.033}_{0.026}$ + 0.12+ 0.07ZH0.78 ± 0.13 ± 0.06 -0.10-0.05 $0.59 \, {}^{+}_{-}\, {}^{0.03}_{0.05}$ + 0.06+ 0.03 $t\bar{t}H + tH$ 0.64 ± 0.12 ± 0.09 ± 0.05 ± 0.08 -0.02

